

Kinetic cross-coupling between non conserved and conserved fields in phase field models

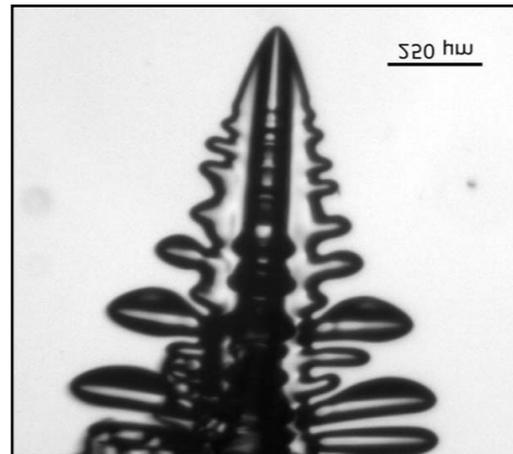
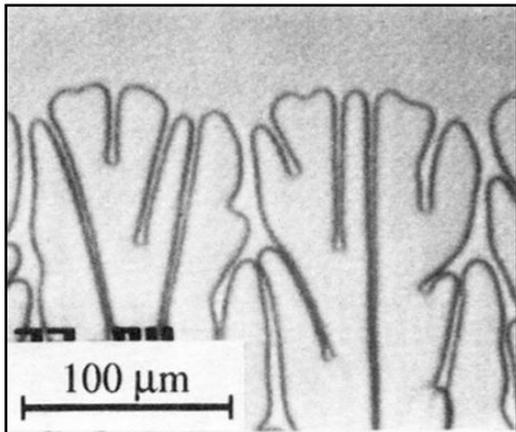
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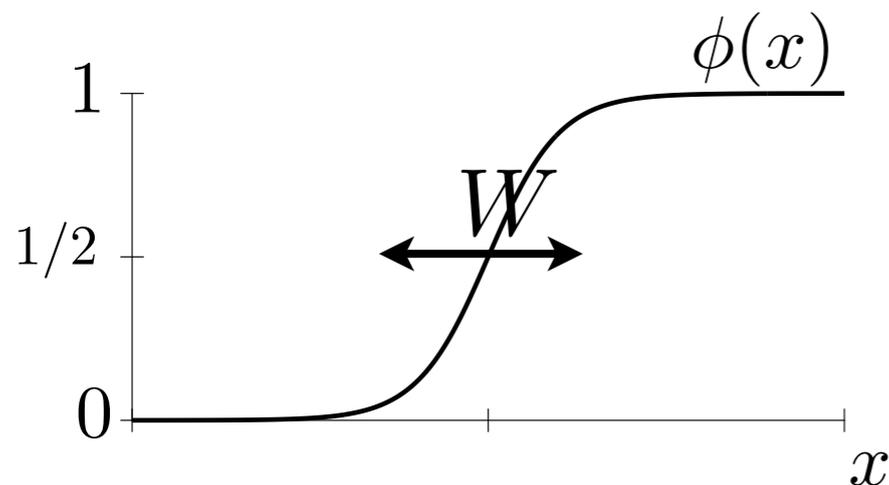
Diffusional growth and kinetic effects



Transport equation in the bulk (diffusion equation)

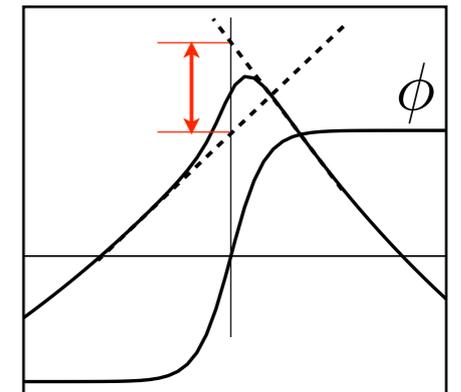
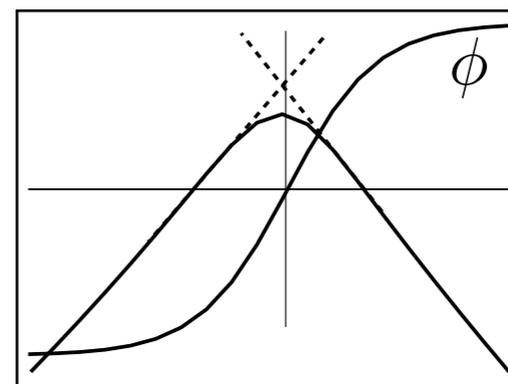
+ boundary conditions at the interfaces :

- energy/mass conservation
- local equilibrium + kinetic effects



thin interface analysis : kinetic effects depend strongly on W

→ which boundary conditions the phase field model corresponds to ?



Phase field models for diffusional growth

→ ϕ : non-conserved field

→ C : conserved field $\dot{C} + \nabla \cdot \mathbf{J} = 0$

$$F = \int_V dV \{ f(\phi) + W^2 |\nabla \phi|^2 + g(\phi, C) \}$$

↑
thermodynamic coupling between phase field and diffusion field

classical phase field model :

$$-\frac{\delta F}{\delta \phi} = \tau \dot{\phi}$$

→ 2 velocity scales : W/τ , D/W

$$-\nabla \frac{\delta F}{\delta C} = \frac{\mathbf{J}}{D(\phi)}$$

Macroscopic approach (sharp-interface description) of diffusional transformations in binary alloys $A_{1-c}B_c$

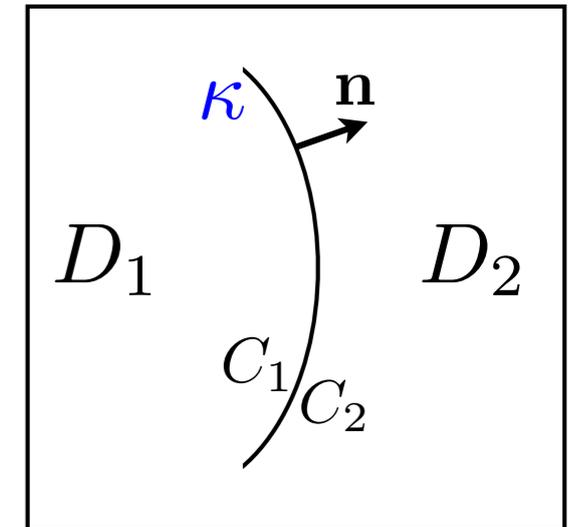
Diffusion equation in the bulk: $\dot{C} = D_i \nabla^2 C$

Conservation of mass at the interface:

$$D_1(\nabla C_1 \cdot \mathbf{n}) + \underbrace{VC_1}_{\text{total number of atoms}} = D_2(\nabla C_2 \cdot \mathbf{n}) + \underbrace{VC_2}_{\text{number of solute atoms}} = J_B$$

total number of atoms

number of solute atoms



Interface kinetics:

$$\Phi_2 - \Phi_1 = \mathcal{A}V + \mathcal{B}J_B + d_0\kappa$$

grand potential:

$$\Phi_i = f_i(C_i) - C_i\mu_i$$

chemical potential:

$$\mu_i = f'_i(C_i)$$

$$\mu_2 - \mu_1 = \mathcal{B}V + \mathcal{C}J_B$$

→ 3 velocity scales describing interface kinetics

Non diagonal model (I)

classical diagonal model :

$$-\frac{\delta \mathcal{F}}{\delta \phi} = \tau \dot{\phi} + 0$$
$$-\nabla \frac{\delta \mathcal{F}}{\delta C} = 0 + \frac{\mathbf{J}}{D(\phi)}$$

non diagonal model :

$$-\frac{\delta \mathcal{F}}{\delta \phi} = \tau \dot{\phi} + (MW \nabla \phi) \cdot \mathbf{J}$$

→ 3 velocity scales :
 W/τ , D/W , $1/M$

$$-\nabla \frac{\delta \mathcal{F}}{\delta C} = (MW \nabla \phi) \dot{\phi} + \frac{\mathbf{J}}{D(\phi)}$$

positiveness of the Onsager matrix: $M^2 < \frac{\tau}{\max[D(\phi)W^2(\nabla \phi)^2]}$

Non diagonal model (2)

determinant > 0



$$\left[1 - \frac{D(\phi)(MW\nabla\phi)^2}{\tau} \right] \dot{\phi} = -\frac{1}{\tau} \frac{\delta\mathcal{F}}{\delta\phi} + \frac{MW D(\phi)}{\tau} \nabla\phi \cdot \nabla \frac{\delta\mathcal{F}}{\delta C}$$

$$\dot{C} = \nabla \cdot \left[D(\phi) \nabla \frac{\delta\mathcal{F}}{\delta C} + MW D(\phi) \dot{\phi} \nabla\phi \right]$$

Non diagonal model (2)

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$$\dot{C} = \nabla \cdot \left[D(\phi) \nabla \frac{\delta\mathcal{F}}{\delta C} + \underbrace{MW D(\phi) \dot{\phi} \nabla\phi}_{\text{same structure as anti-trapping current}} \right]$$

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Non diagonal model (2)

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same structure as anti-trapping current

one-sided model: $D_1 \ll D_2$

~~$$D_1(\nabla C_1 \cdot \mathbf{n}) + VC_1 = J_B = D_2(\nabla C_2 \cdot \mathbf{n}) + VC_2$$~~

$\longrightarrow J_B = VC_1^{eq} \longrightarrow$ no requirement concerning Onsager symmetry

Thin interface limit

$$A = \int_{-\infty}^{\infty} dx [\tau - 2MW C_{eq}(x)] [\phi'_{eq}(x)]^2 + \int_{-\infty}^{\infty} dx \left[\frac{C_{eq}^2(x)}{D(\phi_{eq})} - \frac{(C_1^{eq})^2}{2D_1} - \frac{(C_2^{eq})^2}{2D_2} \right]$$

$$B = \int_{-\infty}^{\infty} dx MW [\phi'_{eq}(x)]^2 - \int_{-\infty}^{\infty} dx \left[\frac{C_{eq}(x)}{D(\phi_{eq})} - \frac{C_1^{eq}}{2D_1} - \frac{C_2^{eq}}{2D_2} \right]$$

$$C = \int_{-\infty}^{\infty} dx \left[\frac{1}{D(\phi_{eq})} - \frac{1}{2D_1} - \frac{1}{2D_2} \right]$$

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Equilibrium profiles :

$$\begin{cases} \phi_{eq}(x) = \tanh\left(\frac{x}{\sqrt{2}W}\right) \\ C_{eq}(x) = \bar{C} + \Delta C \frac{p[\phi_{eq}(x)]}{2} \end{cases}$$

$$\bar{C} = \frac{C_1^{eq} + C_2^{eq}}{2}$$

$$\Delta C = C_1^{eq} - C_2^{eq}$$

Switching function : $p(\phi) = \frac{15}{8} (\phi - 2\phi^3/3 + \phi^5/5)$

Contrast of diffusivity

$$\frac{1}{D(\phi)} = \frac{1}{\bar{D}} + \frac{1}{\Delta D} q(\phi) \quad \text{with} \quad \begin{cases} \frac{1}{\bar{D}} = \frac{1}{2D_1} + \frac{1}{2D_2}; & \frac{1}{\Delta D} = \frac{1}{2D_1} - \frac{1}{2D_2} \\ q(\phi) = -q(-\phi); & q(\pm 1) = \pm 1 \end{cases}$$

$$\mathcal{A} = \frac{\alpha\tau}{W} \left(1 - \beta\Delta C^2 \frac{W^2}{\bar{D}\tau} \right) - 2\bar{C}\mathcal{B}$$

$$\mathcal{B} = \alpha M - \gamma\Delta C \frac{W}{\Delta D}$$

$$\mathcal{C} = 0$$

$$\alpha = W \int_{-\infty}^{\infty} dx [\phi'_{eq}(x)]^2$$

$$\beta = \frac{1}{\alpha} \int_{-\infty}^{\infty} \frac{dx}{4W} [1 - p^2(\phi_{eq})]$$

$$\gamma = \int_{-\infty}^{\infty} \frac{dx}{2W} [p(\phi_{eq})q(\phi_{eq}) - 1]$$

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$$\mathcal{B} = \alpha M - \gamma\Delta C \frac{W}{\Delta D}$$

$$\mathcal{C} = 0$$

Almgren: $\mu_2 - \mu_1 = \mathcal{B}V + \mathcal{C}J_B = 0$

if $\mathcal{B} = 0$ i.e. $M = \frac{\gamma\Delta C}{\alpha} \frac{W}{\Delta D}$

then $\Phi_2 - \Phi_1 = \mathcal{A}V + d_0\kappa$

with $\mathcal{A} = \frac{\alpha\tau}{W} \left(1 - \beta\Delta C^2 \frac{W^2}{\bar{D}\tau} \right)$

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$$\alpha = W \int_{-\infty}^{\infty} dx [\phi'_{eq}(x)]^2$$

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Karma-Rappel: $1/\Delta D = 0$

$\rightarrow \mathcal{B} = \mathcal{C} = 0$ if $M = 0$

Conclusion

- The kinetic cross coupling is necessary to fully describe interface kinetics.
- allows to have the same number of degrees of freedom in phase field model and sharp-interface description
- solves problems:
 - elimination of temperature jump when finite contrast of diffusivity
 - introduce the Ehrlich-Schwoebel effect in MBE
 - tuning the solute trapping effect in alloys
- open question:
 - introduction of kinetic cross coupling for multi-phase systems (treatment of the triple junction ?)

