



Recent advances in Phase-Field Crystal modeling of heterogeneous crystal nucleation

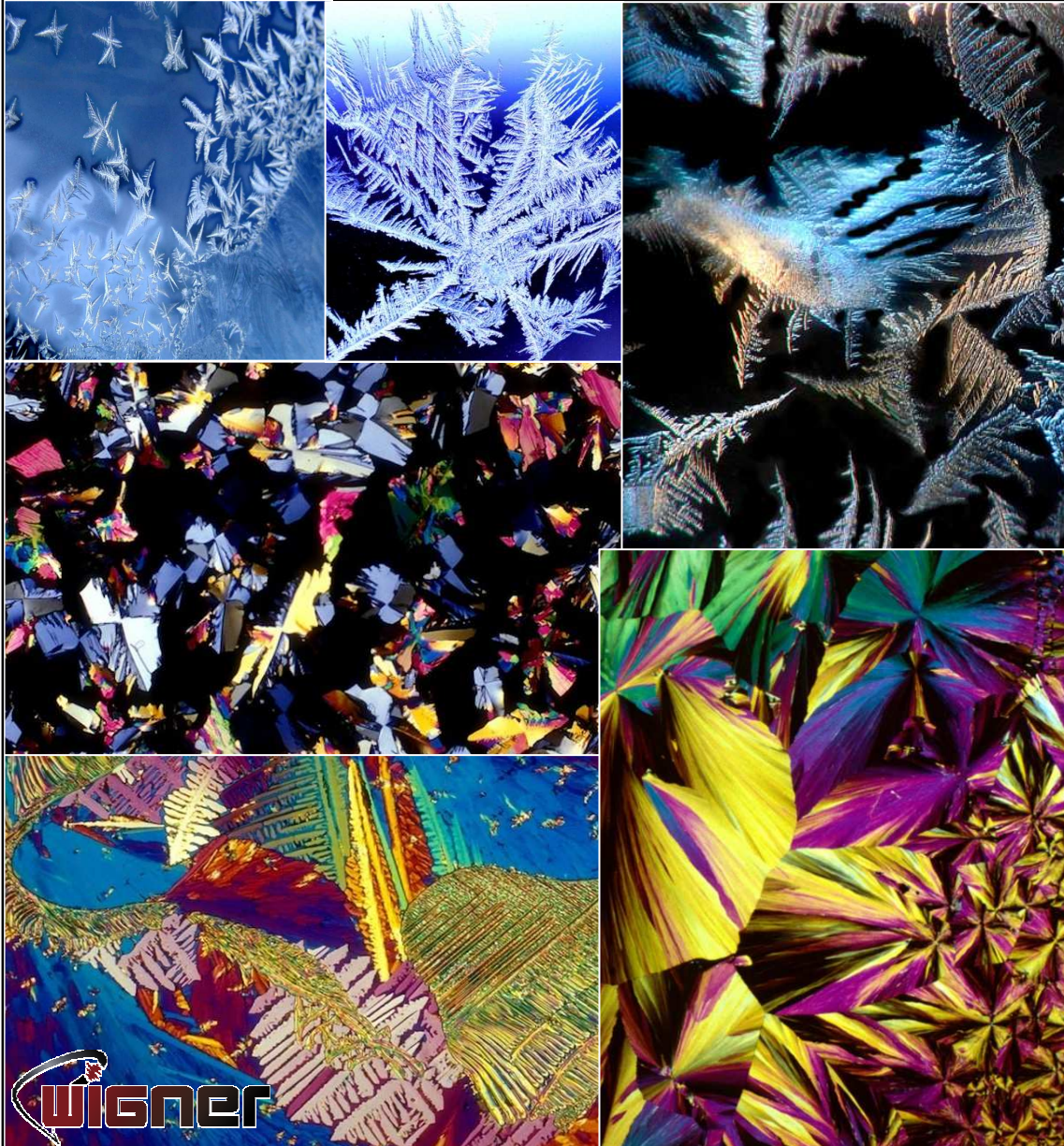
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**Int. Symp. on Phase-field Method – PFM 2014, August 26 – 29, 2014
The Penn Stater Conference Center Hotel, State College, PA, USA**

I. Introduction: Complex polycrystalline structures



Complex patterns evolve
due to the interplay of
nucleation and growth.

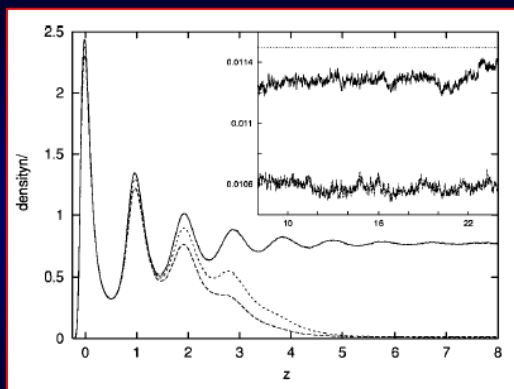
Usually **heterogeneous nucleation**,
during which crystallization is assisted
by foreign walls or particles

Heterogeneous nucleation

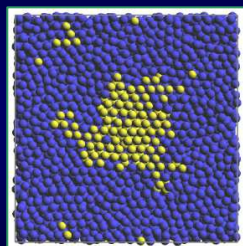
MD/MC Results:

(a) Unstructured walls cause liquid ordering:

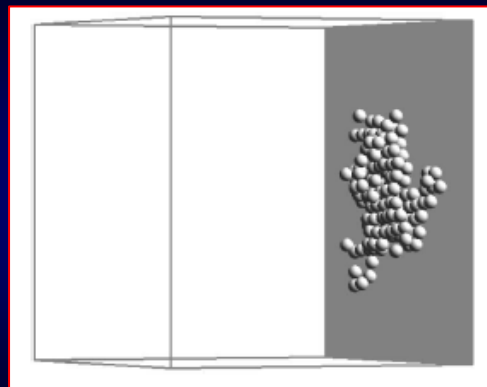
(e.g.: Toxwaerd 2002)



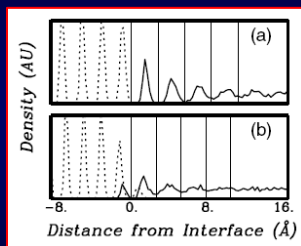
(b) Unstructured wall: HS (111) nearly wets it



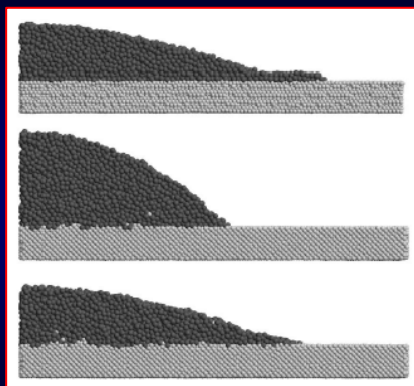
(Auer & Frenkel, PRL 2003)



(c) Liquid ordering at cryst. wall vs orientation:



(Webb & al., PRL 2003)



MD/MC:
the liquid orders at the substrate-liquid interface, which may either help or suppress crystallization depending on the closeness of the ordered liquid structure to that of the crystalline phase.

II. The Phase-Field Crystal Approach: a Simple Dynamical Density Functional Theory

(K. R. Elder *et al.*, PRL, 2002)

The Phase-Field Crystal Model (K.R. Elder et al. PRL 2002)

Classical density functional theory (RY: The free energy is Taylor expanded)

$$\frac{\mathcal{F}}{k_B T} = \int d\vec{r} [\rho(\vec{r}) \ln(\rho(\vec{r})/\rho_\ell) - \delta\rho(\vec{r})] - \frac{1}{2} \int d\vec{r} d\vec{r}' \delta\rho(\vec{r}) C_2(\vec{r}, \vec{r}') \delta\rho(\vec{r}') + \dots,$$

Reference: homogeneous liquid (ℓ)

ρ : time-averaged particle density (number density)

$\delta\rho = \rho - \rho_\ell$

C_2 : two-particle direct correlation function: $C_2(k) = 1 - 1/S(k)$

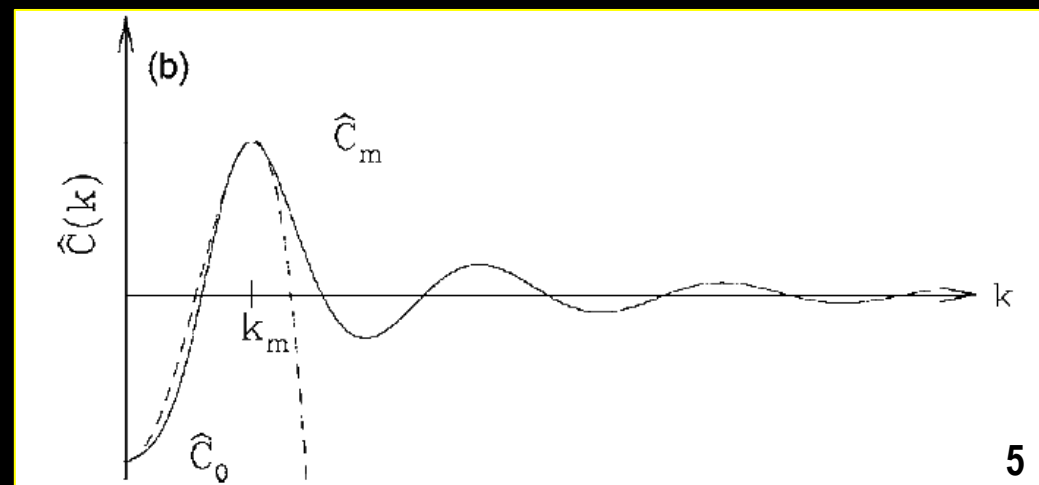
PFC: The two-particle dir. correlation function is Taylor expanded up to 4th order in Fourier space

Physics: in the 3 exp. coeffs. related to the compressibility, bulk modulus and lattice const.

Free energy functional:

Brazovskii/Swift-Hohenberg form:

$$\Delta\tilde{F} = \int d\vec{r} \left\{ \frac{\psi}{2} [r + (1 + \tilde{\nabla}^2)^2] \psi + \frac{\psi^4}{4} \right\}$$

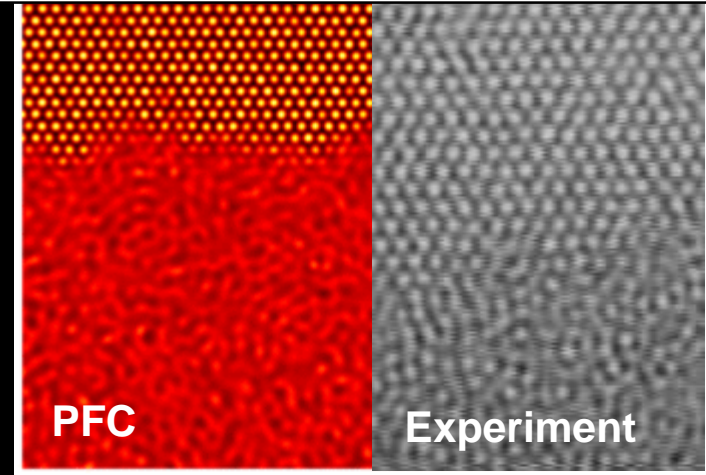


Euler-Lagrange eq. (ELE)
(extremes of free energy)

$$\left. \frac{\delta \Delta \tilde{F}}{\delta \psi} = \frac{\delta \Delta \tilde{F}}{\delta \psi} \right|_{\psi_0}$$

Equilibrium properties:

- phase diagram
- interface free energies
- nucleation barrier, etc.

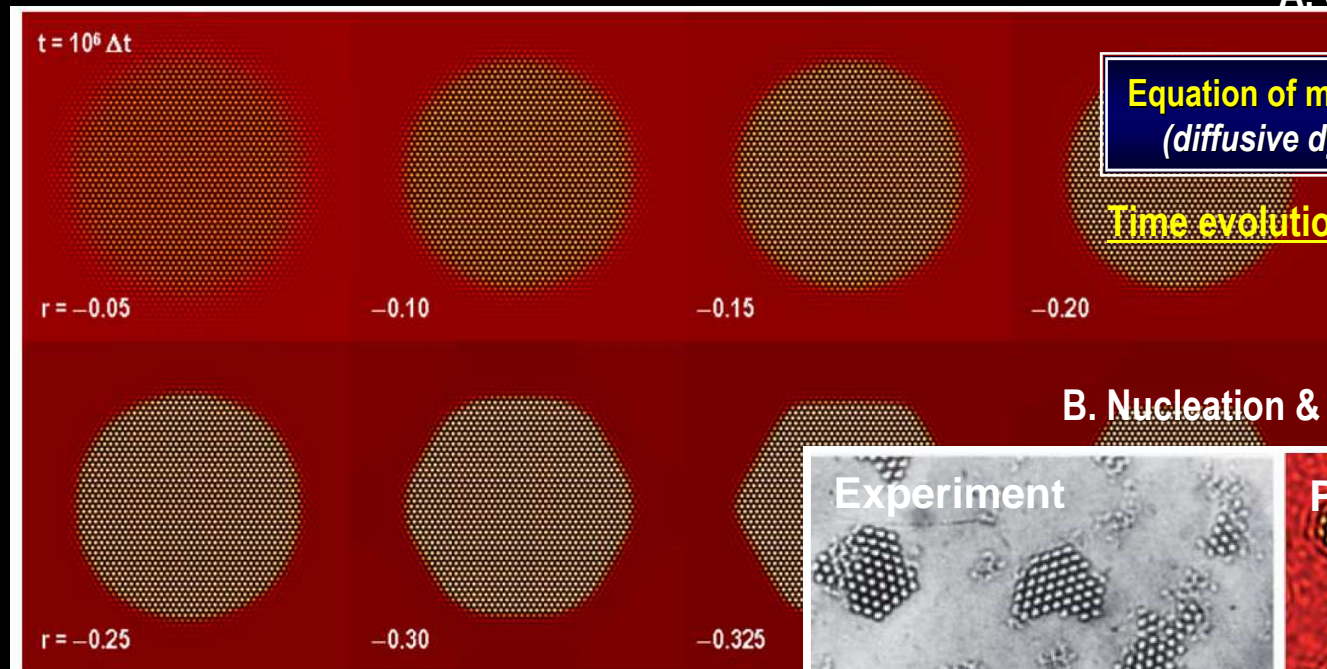


PFC

Experiment

A. Crystal-liquid interface

Equilibrium shapes. Parameter r regulates anisotropy.

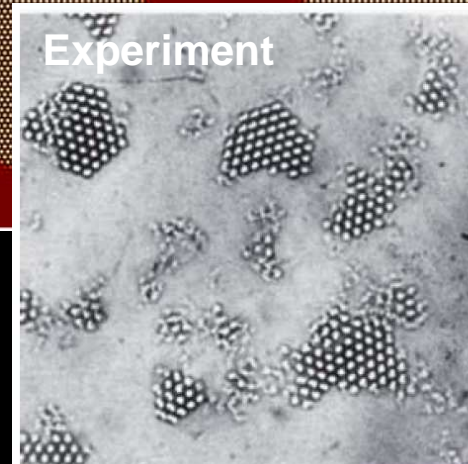


Equation of motion (EOM):
(diffusive dynamics - colloids)

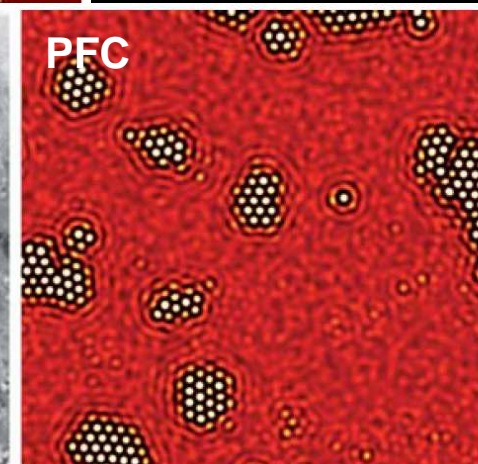
Time evolution:

$$\frac{\partial \psi}{\partial \tilde{t}} = \tilde{\nabla} \left\{ \tilde{\nabla} \frac{\delta \Delta \tilde{F}}{\delta \psi} + \zeta \right\}$$

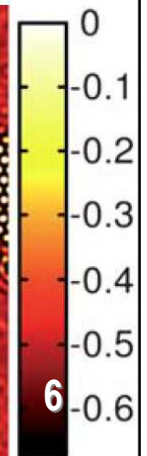
B. Nucleation & growth in colloids



Experiment

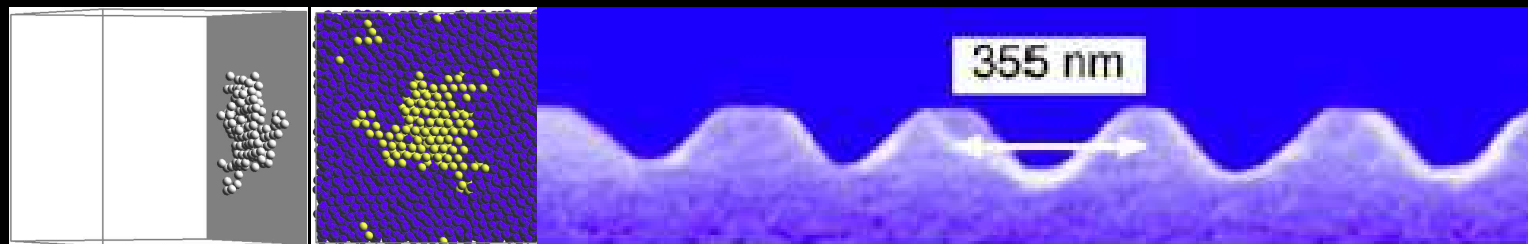


PFC

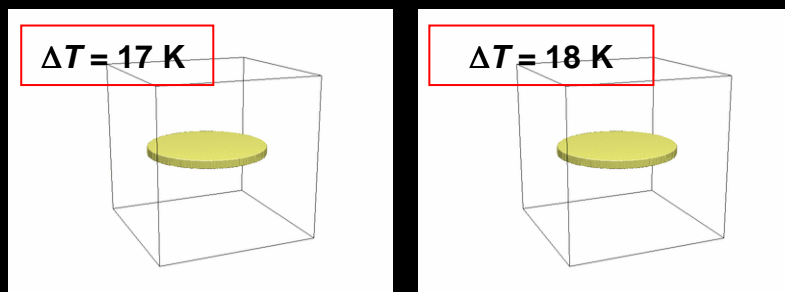


III. PFC Results for Heterogeneous Nucleation

A. Crystallization induced by flat & modulated surfaces?



B. Particle-induced freezing?

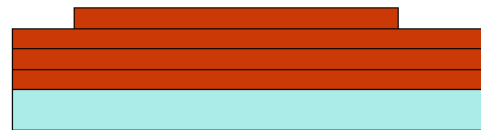


A. Crystallization on flat/modulated surfaces in 2D

A.1: Crystallization on flat walls:

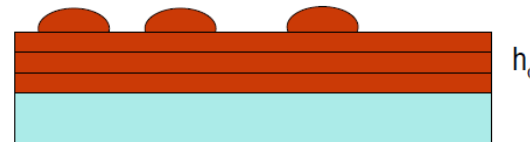
Three Growth Modes in Heteroepitaxy (A/B) (Bauer, 1958)

(a) **2D or layer-by-layer**
(Frank - Van der Merwe)



$$\Delta\gamma = \gamma_o + \gamma_i - \gamma_s < 0$$

(b) **2D followed by 3D
islanding**
(Stranski - Krastanov)



$$\Delta\gamma \sim 0$$

(c) **3D islanding**
(Volmer - Weber)



$$\Delta\gamma > 0$$

γ_o – free energy of **crystal – liquid** interface
 γ_i – free energy of **crystal – substrate** interface
 γ_s – free energy of **substrate – liquid** interface

Substrate in PFC:

Extra term added to free energy density:

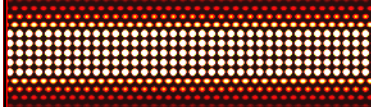
$$f_{\text{PFC}} + V(\underline{r}) \psi(\underline{r})$$

Potential: $V(\mathbf{r}) = [V_0 - V_1 S(a_s, \mathbf{r})] h(\mathbf{r})$

V_0 and V_1 – constants
 $S(a_s, \underline{r})$ – SM approximant
 $h(\underline{r})$ – envelop function

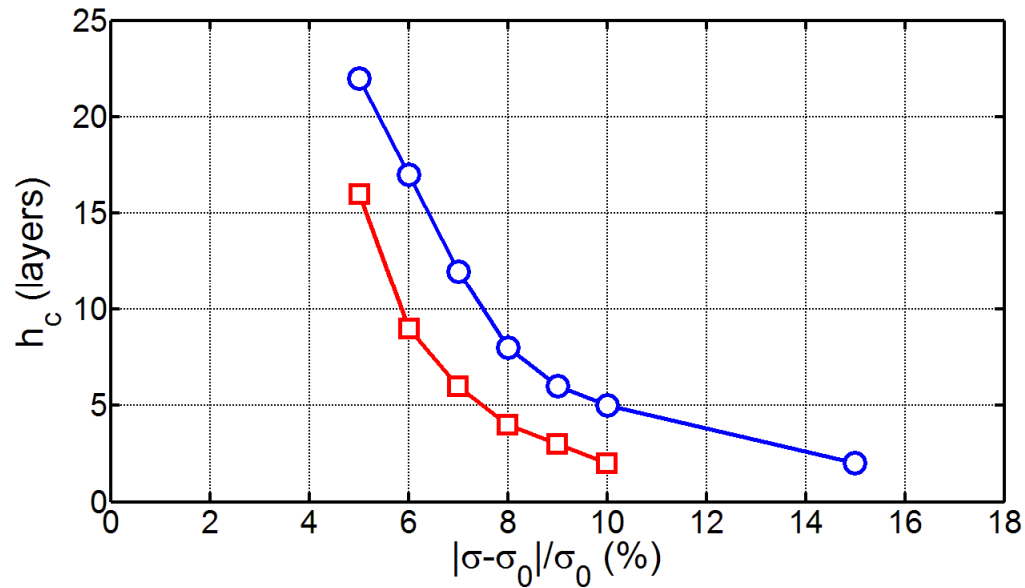
(a) Crystallization on flat walls in cooling (EOM): $\psi_0 = -0.25$; $\varepsilon = 0.1 - 0.1875$ (diffuse SL interface)

$a_s/\sigma = 1.00$



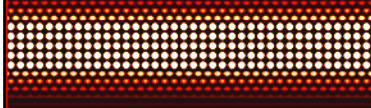
Layer-by-layer
(Frank - van der Meerwe)

$\varepsilon \approx 0.1 - 0.2$
metal-like

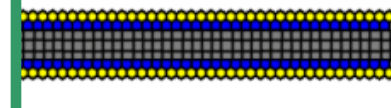
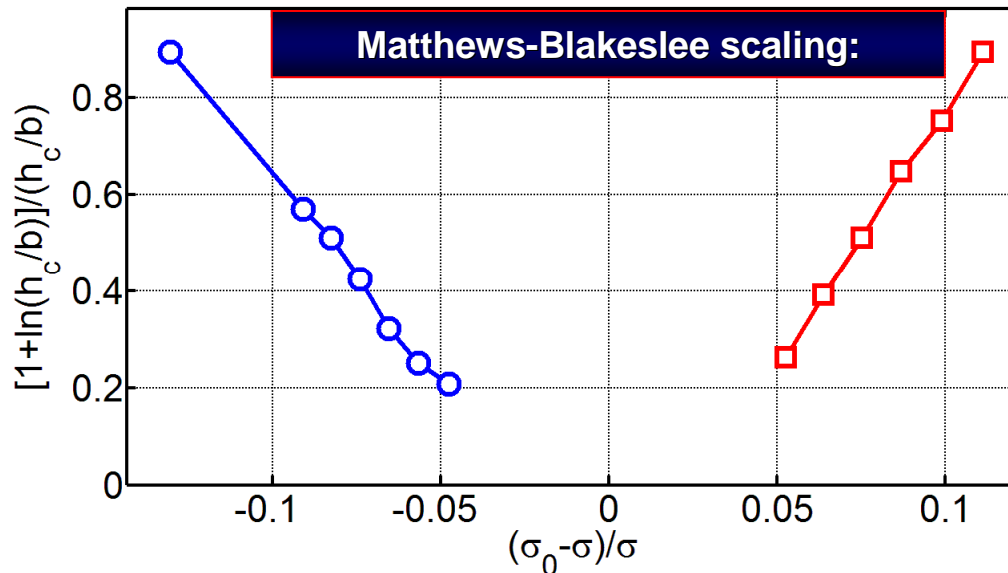


ue; 6 : yellow; 7 : red

$a_s/\sigma = 1.07$

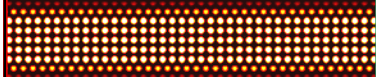


Layer-by-layer follow
(Stranski-Krastanov, Asarc



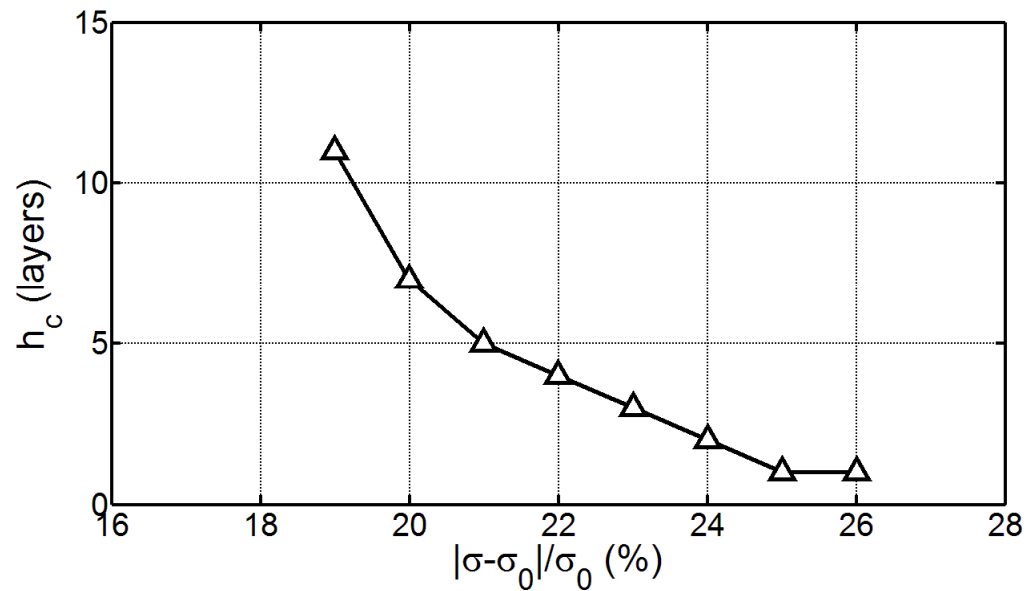
(b) Crystallization on flat walls in cooling (EOM): $\psi_0 = -0.5$; $\varepsilon = 0.4 - 0.75$ (faceted SL interface)

$a_s/\sigma = 1.00$



Layer-by-layer
(Frank - van der Meerwe)

$\varepsilon \approx 0.4 - 0.8$
covalent

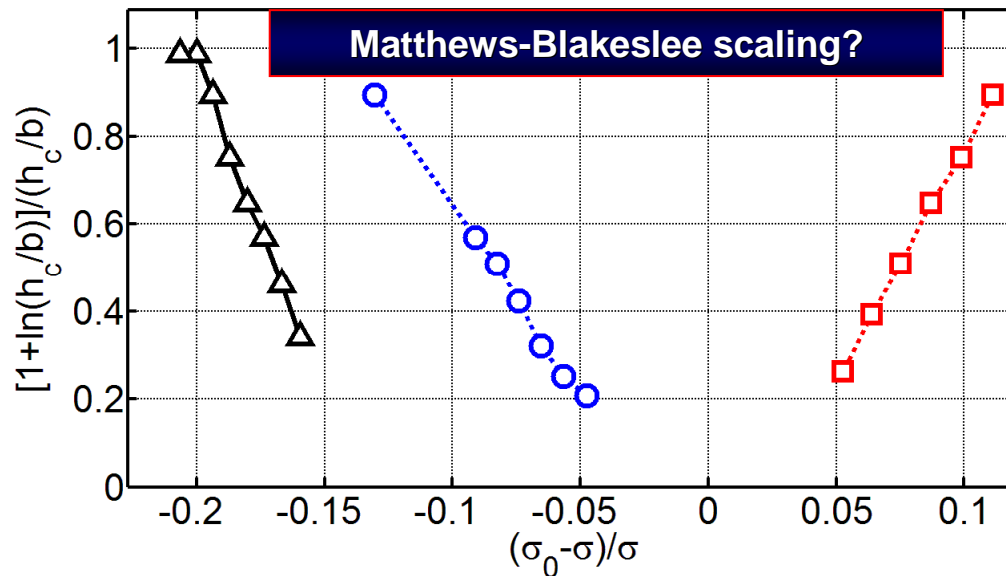


ue; 6 : yellow; 7 : red

$a_s/\sigma = 1.19$



Layer-by-layer follow
(Stranski-Krastanov, Asarc



Crystallization on flat walls in cooling (EOM): $\psi_0 = -0.5$; $\varepsilon = 0.4 - 0.75$

$a_s/\sigma = 1.2$ EOM $\zeta = 0.0$

$a_s/\sigma = 1.2$ EOM $\zeta = 1.0$

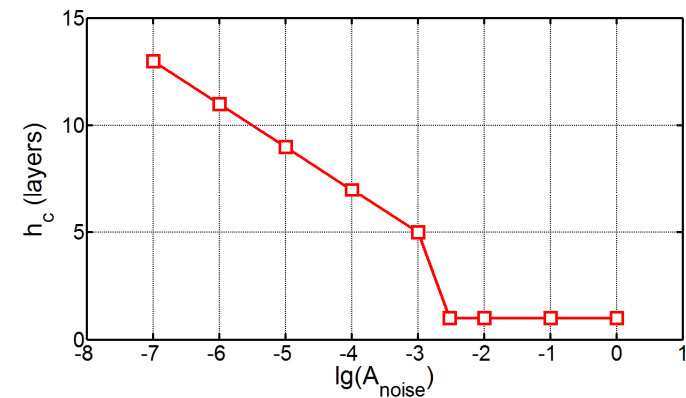
Observations:

Without noise:

- First pseudomorphic solid forms;
- Beyond critical thickness (h_c) misfit dislocations form that release elastic strain energy.

With noise:

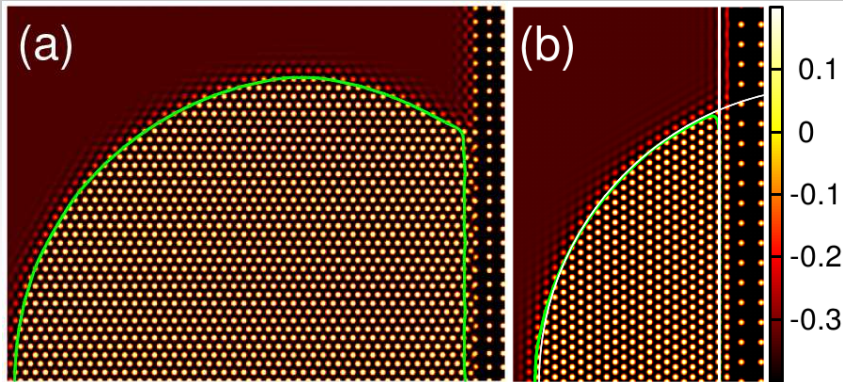
- Periodic defect structure in the first layer from the beginning



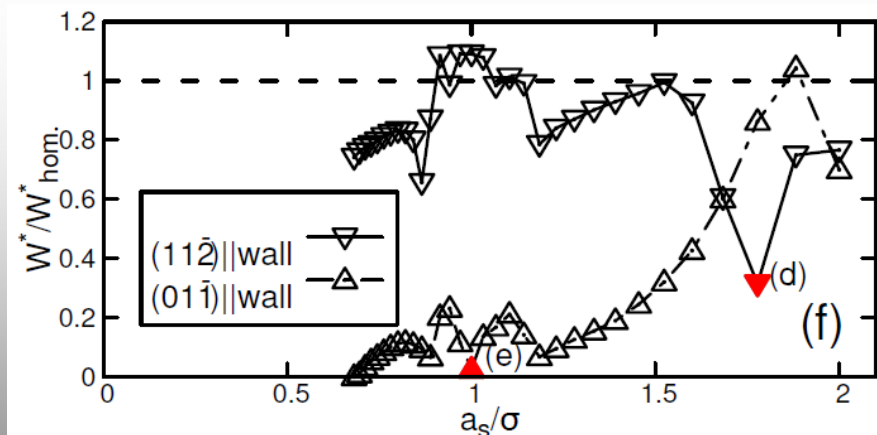
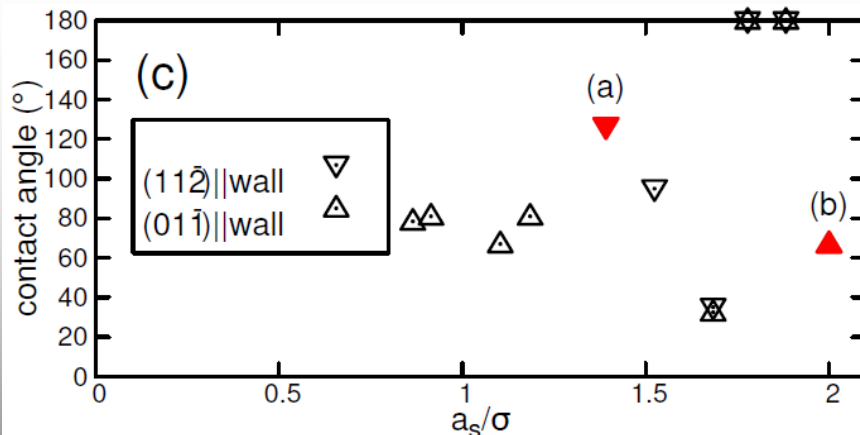
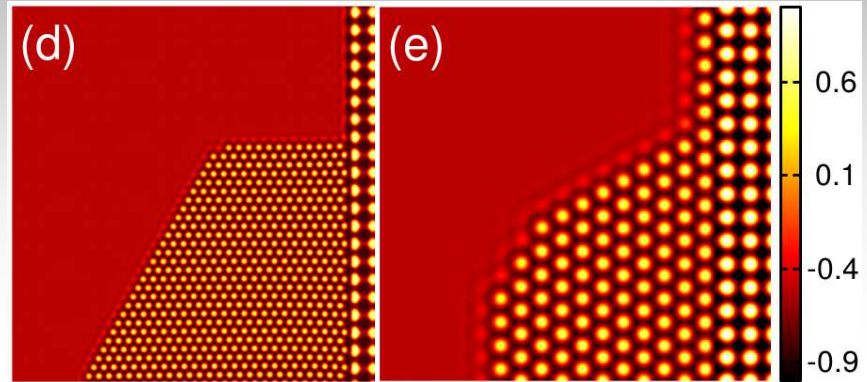
A.2: Island formation (heterogeneous nucleation) on flat interfaces (ELE): (Volmer-Weber)

Potential : $V(\mathbf{r}) = [V_0 - V_1 S(a_s, \mathbf{r})] h(\mathbf{r})$

Nearly isotropic ($r = -0.25$) ; $\psi_0 = -0.341$, $V_0 = V_1 = 0.5$



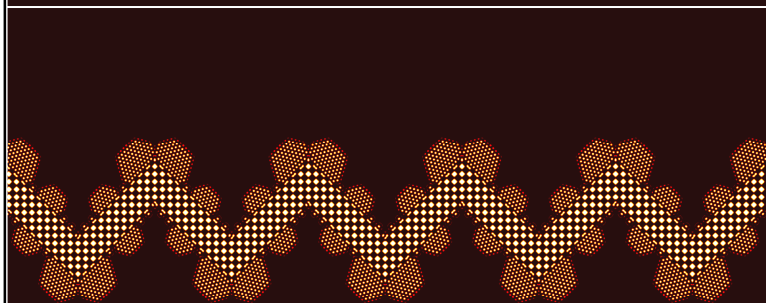
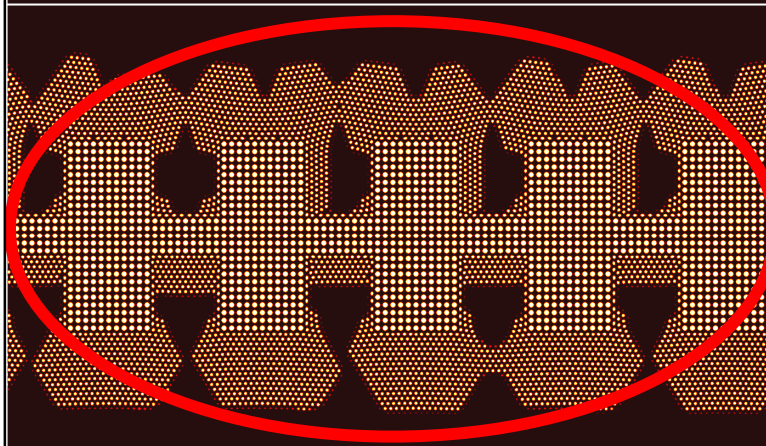
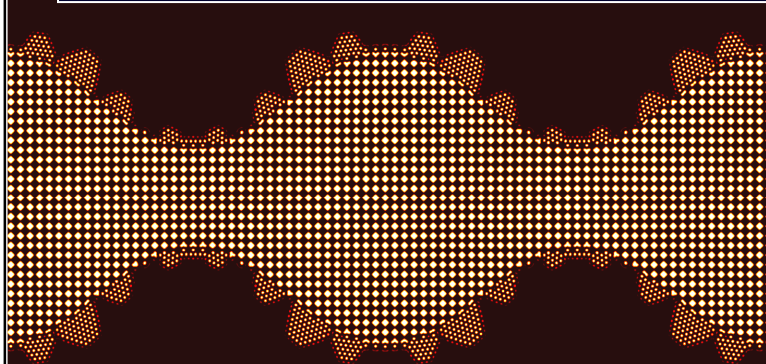
Faceted ($r = -0.5$) ; $\psi_0 = -0.51415$, $V_0 = 0$, $V_1 = 0.65$



Tóth et al. PRL (2012).

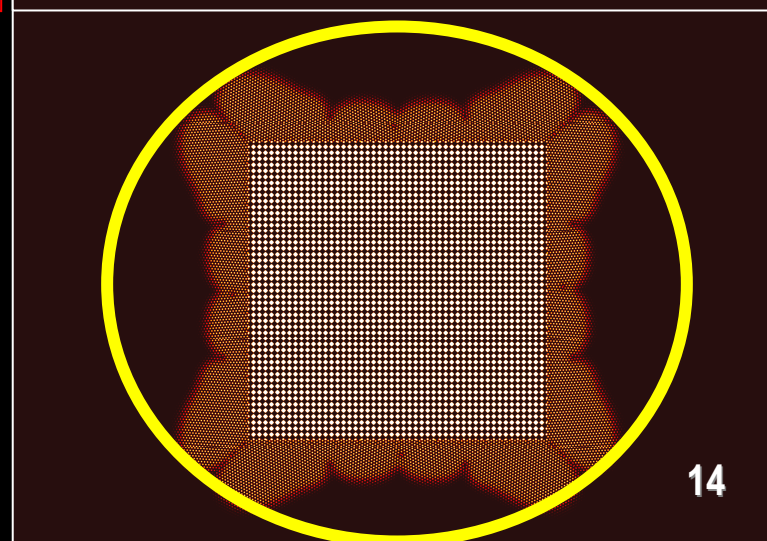
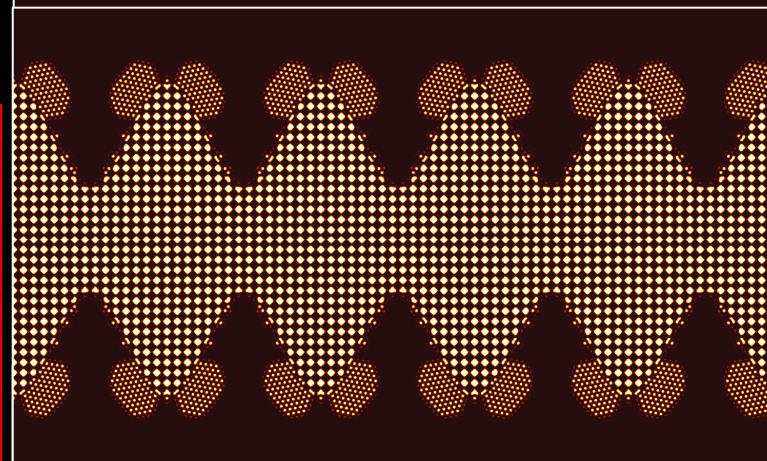
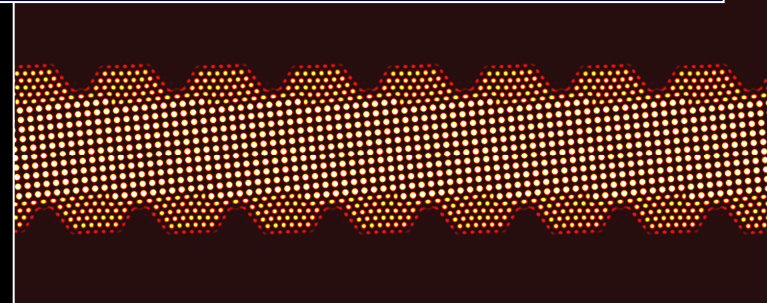
Contact angle (ϑ) and barrier height (W^*)
are nonmonotonic functions of a_s

A.3: Heteroepitaxy on various substrate configurations: ($a_s/\sigma = 1.2$; $\psi_0 = -0.5$; ; $\epsilon = 0.4 - 0.75$)



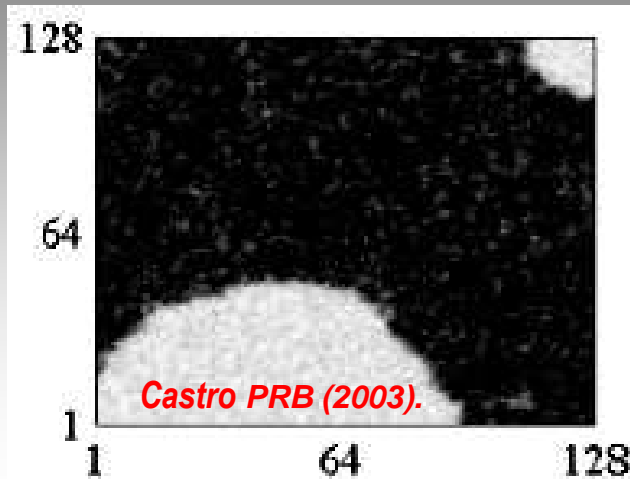
Various shapes:

- sinus & rectangular grooves,
- saw-tooth,
- atomic-scale steps,
- square, etc.



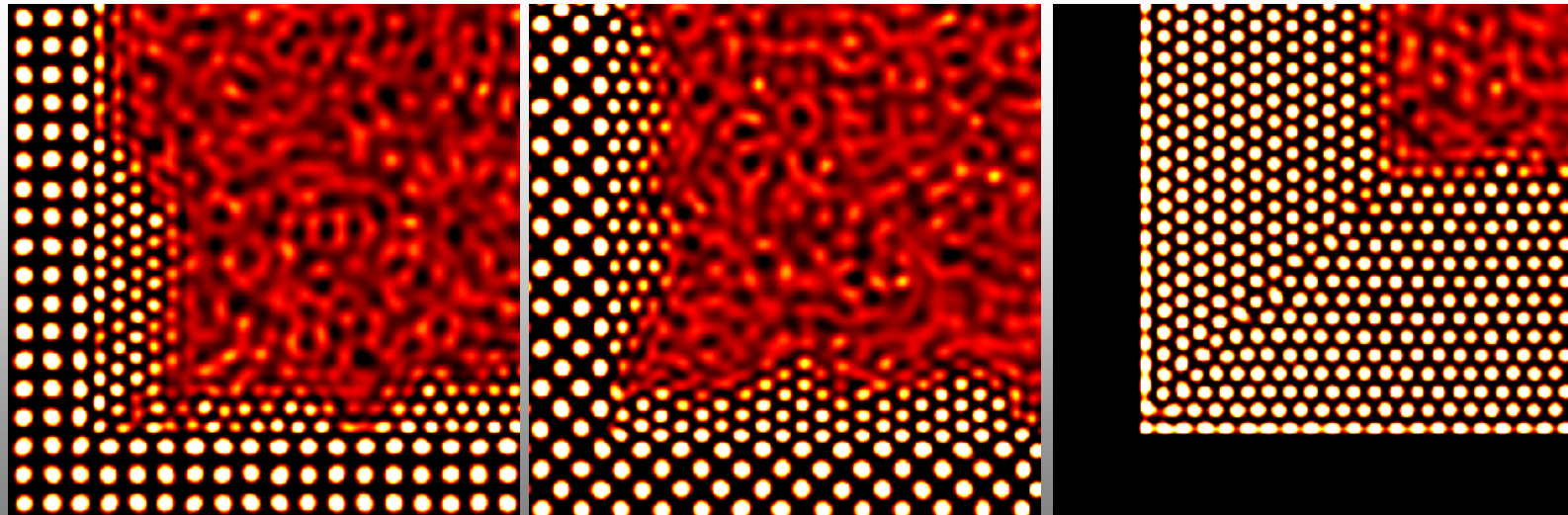
A.4: Heterogeneous nucleation in rectangular corner (EOM):

(a) Conventional PF & CNT:
The corner is preferable.



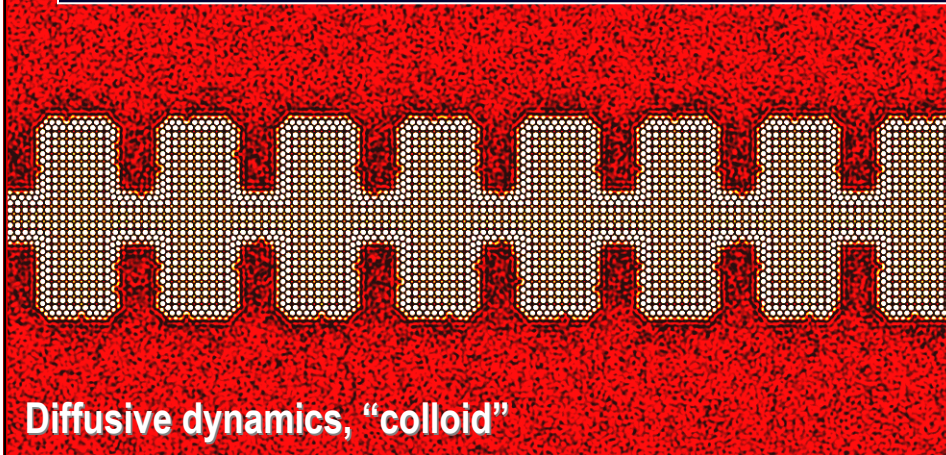
(b) PFC: This corner is NOT preferable!!!

$r = -0.25$
 $\psi_0 = -0.32$
 $\alpha = 0.1$
 $a_s/\sigma = 1.39$

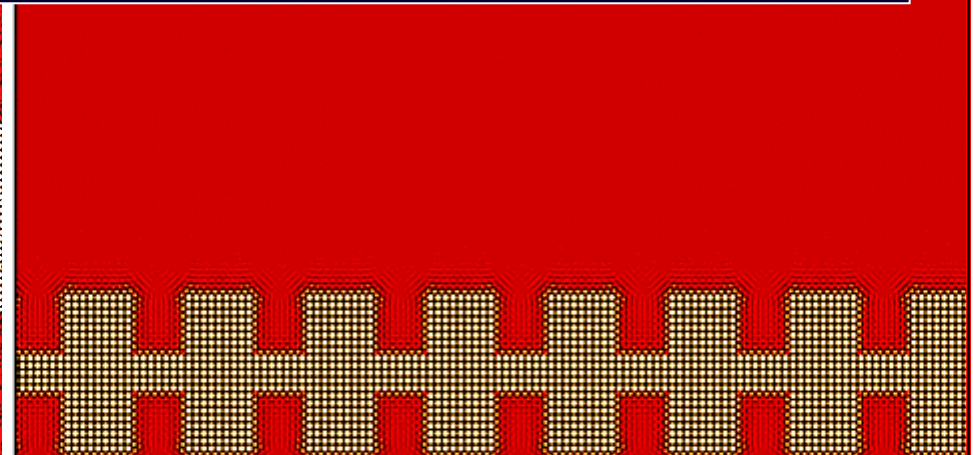
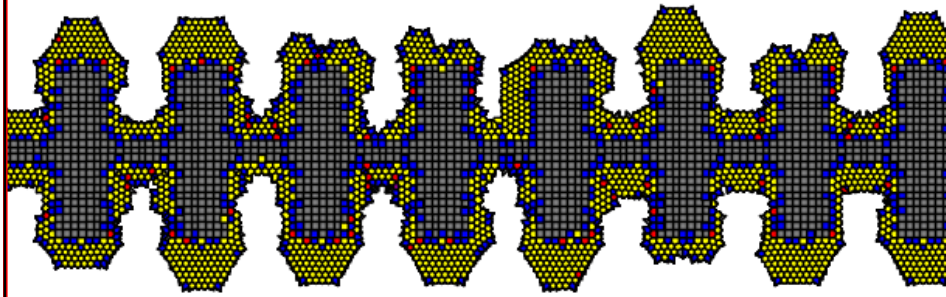


Gránásy et al. Philos. Mag. (2011).

A.5: Rectangular grooves: Isothermal freezing of “colloidal” and “metallic” liquids in PFC:



Diffusive dynamics, “colloid”



Hydrodynamic theory, “Fe”

Observations:

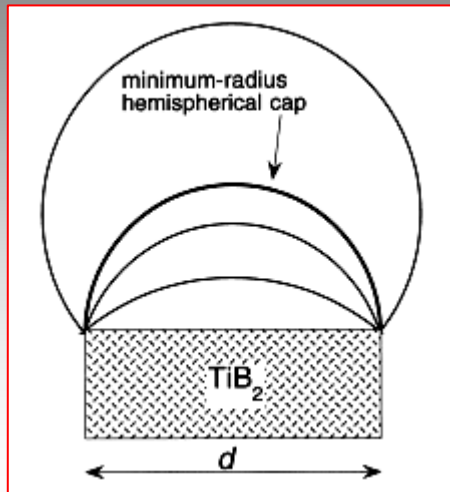
I. Diffusive dynamics:

- Depletion zone
- Faceted crystals all the time

II. Hydrodynamic theory:

- No depletion zone
- Small anisotropy
- Capillary waves

B. Particle-induced freezing in 2D & 3D:



(Greer et al., Acta Mater., 2002)

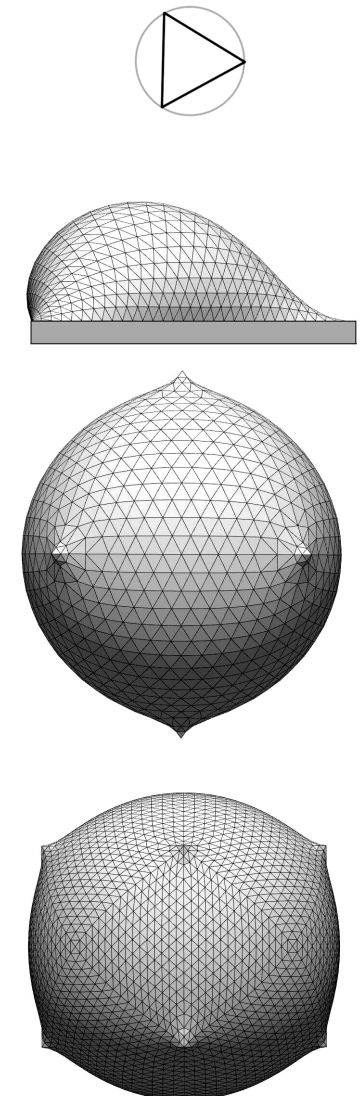
- Cylindrical particles ~ wet by the crystal on top/bottom, not on sides;

(e.g., Al + Al-Ti-B inoculant \rightarrow Ti_2B particles with $AlTi_3$ coating on {0001} faces different contact angles on different faces)

- Free growth for

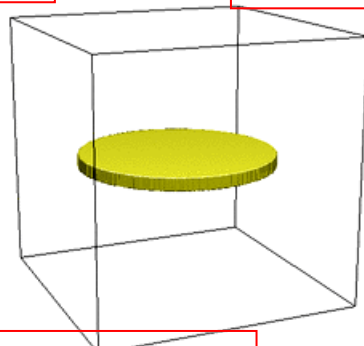
$$\Delta T > \Delta T_c = \frac{4\gamma_{SL}T_m}{Ld}$$

- PFT simulations $\rightarrow \Delta T_c \propto 1/d$; $\Delta T_c <$ classical



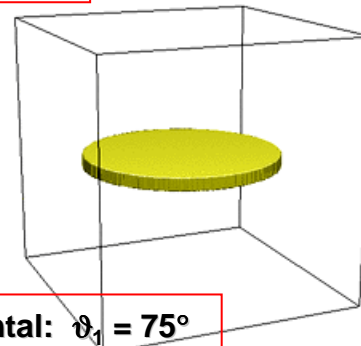
$\Delta T = 17$ K

$d = 30$ nm



40 nm \times 40 nm \times 40 nm

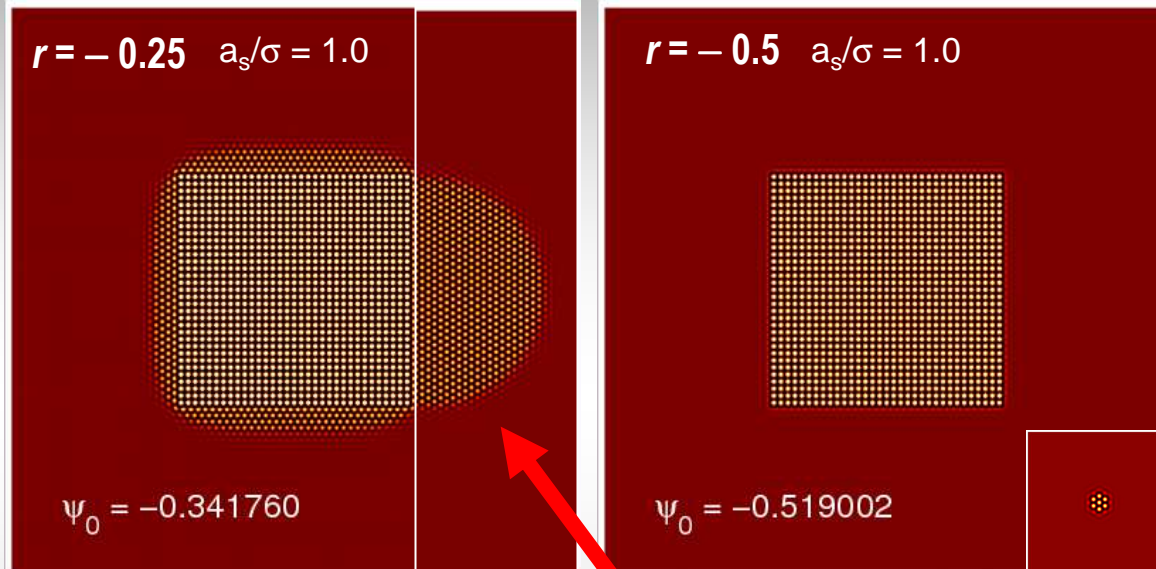
$\Delta T = 18$ K



Horizontal: $\theta_1 = 75^\circ$
Vertical: $\theta_2 = 175^\circ$

B.1 Free-growth limited mode of particle induced freezing in 2D (ELE)

EL solutions for increasing driving force:



Results:

- Small anisotropy: Greer's model OK
- Faceted: free-growth at a much larger driving force

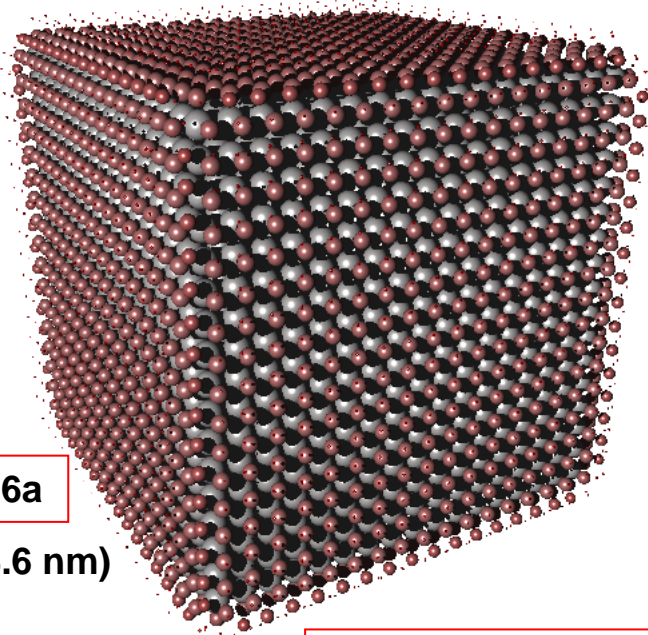
Tóth et al. PRL (2012)

Homogeneous nuclei at
the critical driving force

B.2 Free-growth limited mode of particle induced freezing in 3D (ELE)

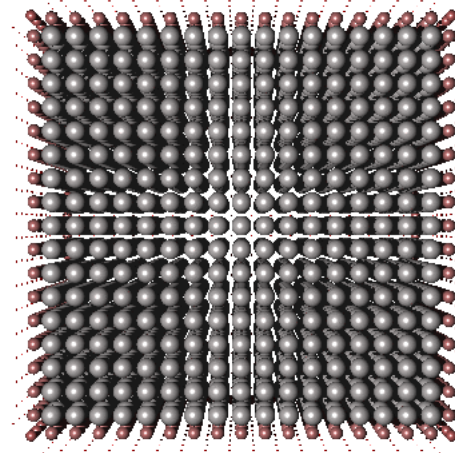
$$r = -0.25$$

SC substrate
Cubic shape
 $a_s/a_{\text{BCC}} = 1$



$$L = 16a$$

($L \sim 4.6 \text{ nm}$)

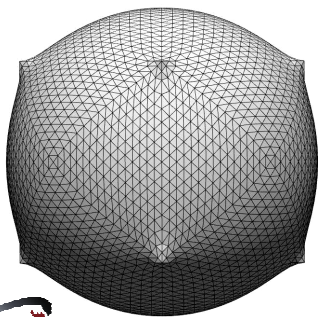


Tóth et al. PRL (2012)

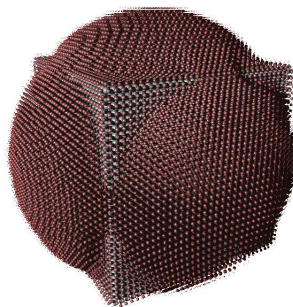
$$\psi_0 = -0.3540$$

256 × 256 × 256 grid

$$\psi_0 = -0.3540$$

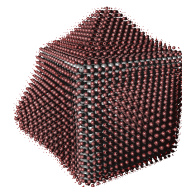


($L \sim 9.2 \text{ nm}$)



$L = 32a$

($L \sim 4.6 \text{ nm}$)



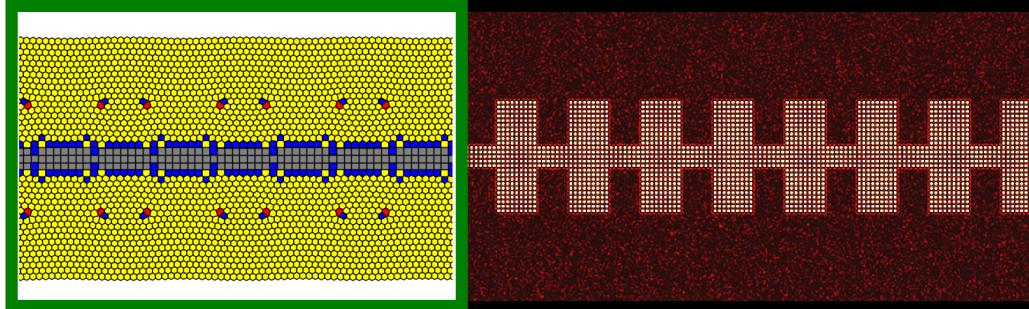
$L = 16a$

The critical shape preceding free growth depends on size!

IV. Summary:

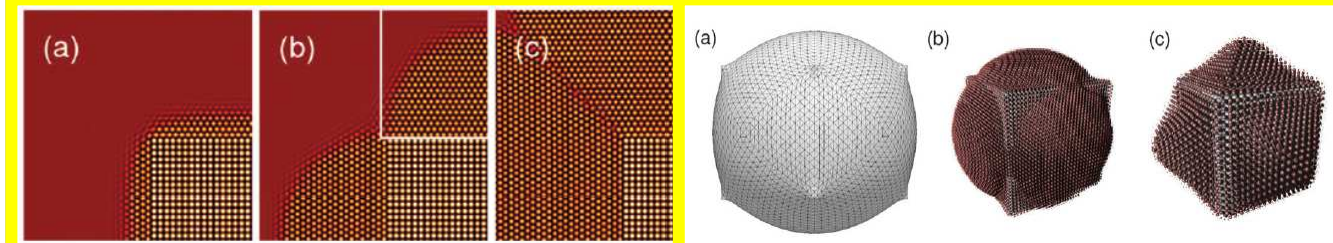
A. Nucleation on walls:

- Epitaxial growth for $a_s = \sigma$
- Surface spinodal for $a_s \neq \sigma$
- Three modes of heteroepitaxy (Asaro-Tiller-Grinfeld instability)



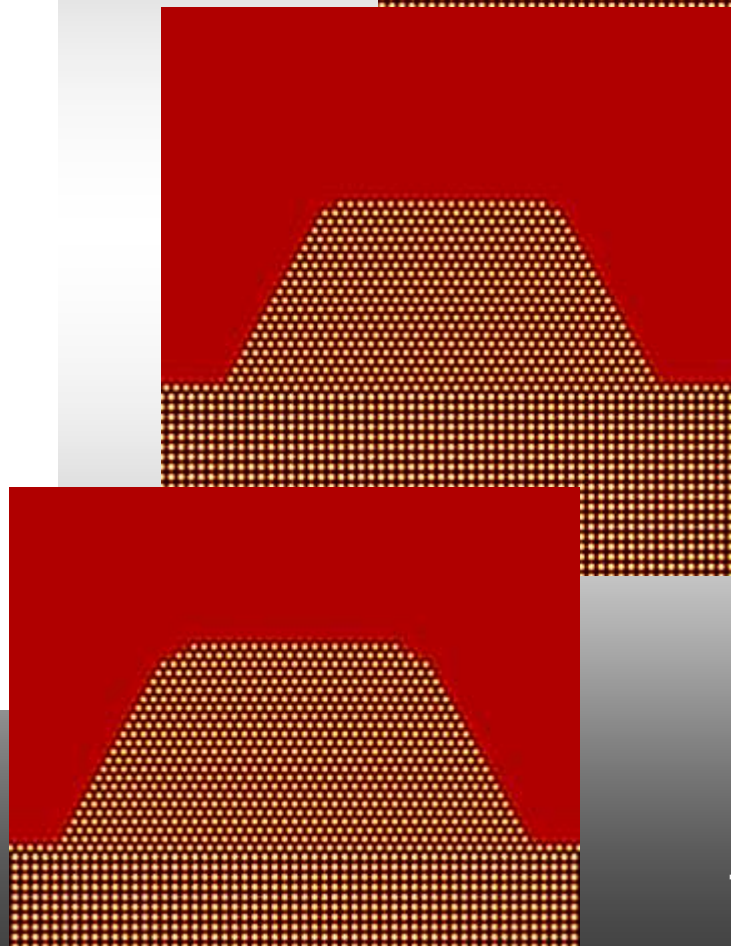
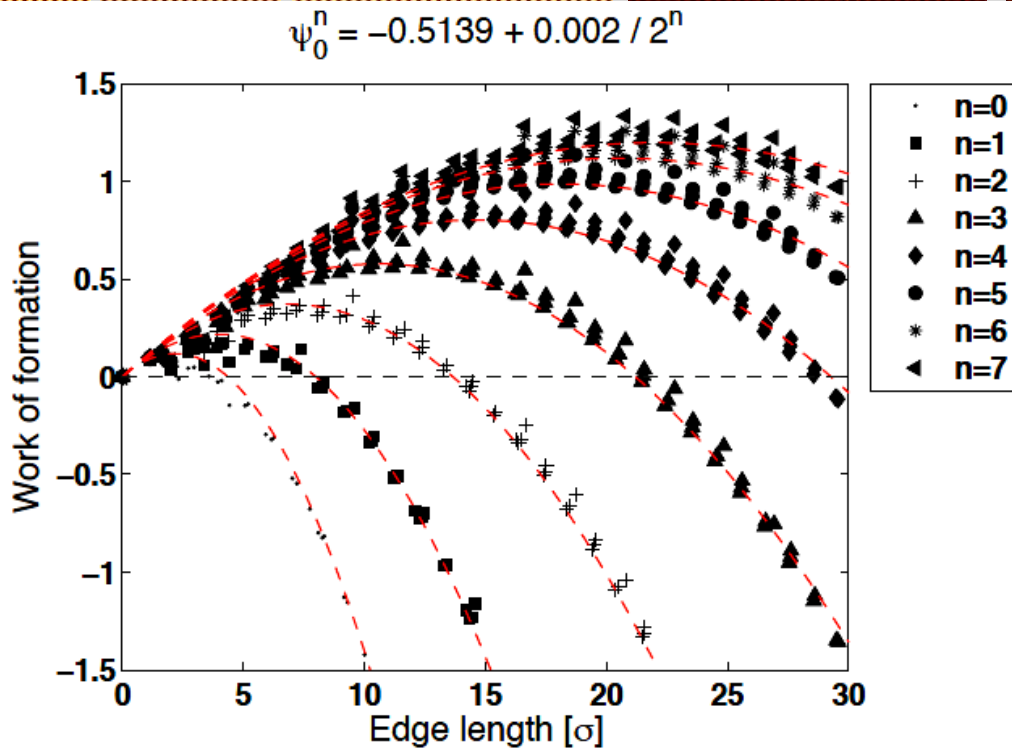
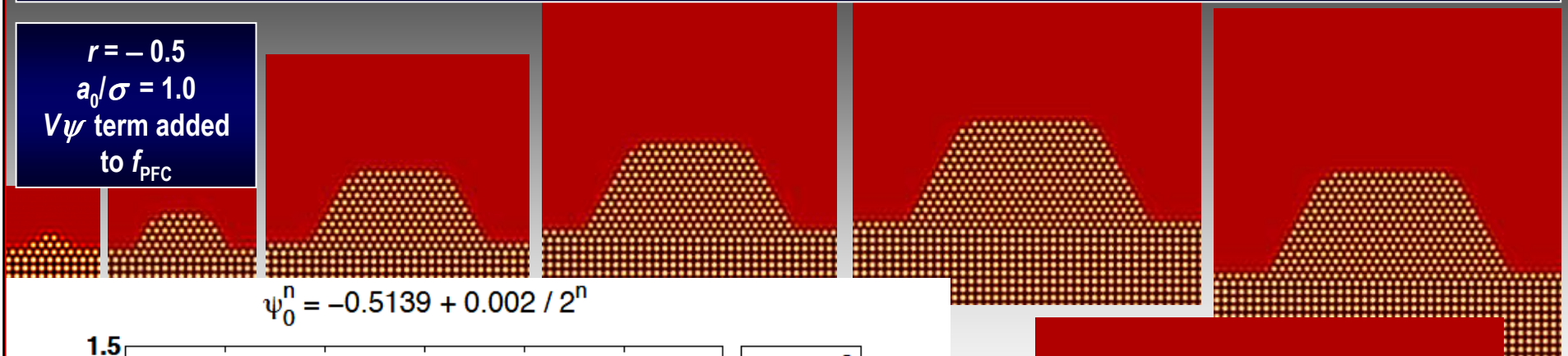
B. Particle induced freezing:

- Works for small ϵ
- Does not work for large ϵ
- Size dependent critical shape
- Non-monotonic dependence of barrier & contact angle on lattice mismatch



Heterogeneous crystal nuclei in 2D (EL equation): (Tóth *et al.*, J. Phys.: Condens. Matter 22, art. no. 364101 (2010))

$r = -0.5$
 $a_0/\sigma = 1.0$
 $V\psi$ term added
 to f_{PFC}



Results:

- Adsorbed monolayer !
- Contact angle: 60 °, determined by crystal structure

Hydrodynamic theory of freezing (Tóth *et al.*, J. Phys.: Condens. Matter 27, art. no. 055001 (2014))

Fluctuating Nonlinear Hydrodynamics

Following L. D. Landau & E. M. Lifshitz (1959)

$$\begin{aligned} \text{Navier-Stokes:} \quad & \frac{\partial \mathbf{p}}{\partial t} + \nabla \cdot (\mathbf{v} \otimes \mathbf{p}) = \nabla \cdot [\mathbb{R}(\rho) + \mathbb{D}(\mathbf{v}) + \mathbb{S}] \\ \text{continuity:} \quad & \frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{p} = 0 \end{aligned}$$

- Variables: $\rho(\mathbf{r}, t)$ (mass density) and $\mathbf{p}(\mathbf{r}, t)$ (momentum density), $\mathbf{v} = \mathbf{p}/\rho$
- $\mathbb{R}(\rho)$: reversible stress tensor (to be defined on thermodynamical basis)
- $\mathbb{D}(\mathbf{v})$: dissipative stress tensor, $\mathbb{D} = \mu_S [(\nabla \otimes \mathbf{v}) + (\nabla \otimes \mathbf{v})^T] + \mu_B (\nabla \cdot \mathbf{v})$
- \mathbb{S} : Stochastic momentum noise with the correlator^[2]:

$$\langle S_{ij}(\mathbf{r}, t) S_{kl}(\mathbf{r}', t') \rangle = (2k_B T \mu_S) \left[(\delta_{ik} \delta_{jl} + \delta_{jk} \delta_{il}) + \left(\frac{\mu_B}{\mu_S} - \frac{2}{3} \right) \delta_{ij} \delta_{kl} \right] \delta_{\mathbf{r}, \mathbf{r}'} \delta_{t, t'}$$

Problems related to applying $F[\rho]$ of a CDFT of freezing in the Navier-Stokes equation

Main problem: The crystal is represented by lattice periodic field distribution!
 $\Rightarrow \mathbf{v} = \mathbf{p}/\rho$ is **singular** & generates **spurious interatomic flows** - UNPHYSICAL

From the viewpoint of *scales*...

The validity limit of the Navier-Stokes is approx. 100× effective molecular diameter...
 (... below which the definition of **the continuum velocity field** fails.)

Scale separation

Coarse-graining (CG):

$$\hat{\mathbf{v}} = \hat{\mathbf{p}}/\hat{\rho}$$

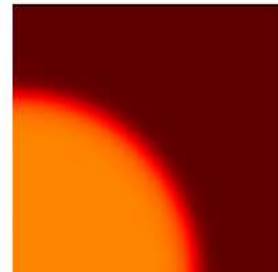
Apply $\nabla(\hat{\mathbf{v}} \otimes \mathbf{p})$ and $\mathbb{D}(\hat{\mathbf{v}})$
in the Navier-Stokes

and

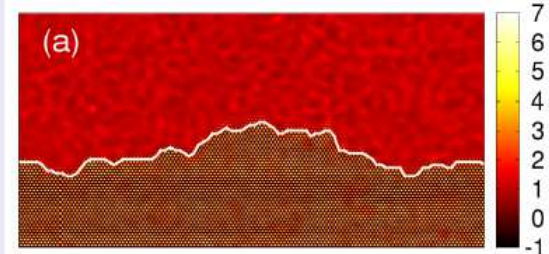
$$\nabla \cdot \mathbb{R} = -\rho_0 \cdot \nabla \frac{\delta F[\rho]}{\delta \rho}$$



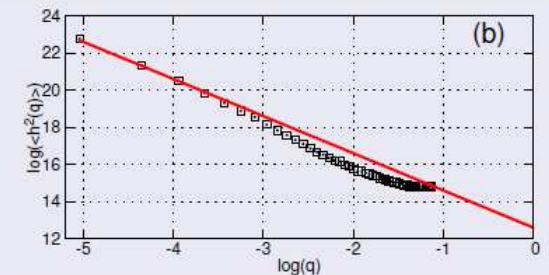
CG \rightarrow



Fluctuating crystal-liquid front



Capillary wave spectrum



Crystalline anisotropy

CW theory: $\langle |h(q)|^2 \rangle = (k_B T)/(q^2 \tilde{\gamma} L)$

Stiffness: $\tilde{\gamma} = \gamma(\theta) + \gamma''(\theta)$

For 6fold symmetry the estimation is

$$\epsilon \approx |(\tilde{\gamma}_{max} - \tilde{\gamma}_{min})/(\tilde{\gamma}_{max} + \tilde{\gamma}_{min})|/35$$

$\Rightarrow \epsilon \lesssim 0.002$ (reasonable in PFC)

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