

1 **A climate sensitivity estimate using Bayesian fusion**  
2 **of instrumental observations and an Earth System**  
3 **model**

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15 **Abstract.** Current climate model projections are uncertain. This uncer-  
16 tainty is partly driven by the uncertainty in key model parameters such as  
17 climate sensitivity ( $CS$ ), vertical ocean diffusivity ( $K_v$ ), and strength of an-  
18 thropogenic sulfate aerosol forcing. These parameters are commonly estimated  
19 using ensembles of model runs constrained by observations. Here we obtain  
20 a probability density function (pdf) of these parameters using the Univer-  
21 sity of Victoria Earth System Climate Model (UVic ESCM) - an interme-  
22 diate complexity model with a dynamic three-dimensional ocean.

23 Specifically, we run an ensemble of UVic ESCM runs varying parameters  
24 that affect  $CS$ , ocean vertical diffusion, and the effects of anthropogenic sul-  
25 fate aerosols. We use a statistical emulator that interpolates the UVic ESCM  
26 output to parameter settings where the model was not evaluated. We adopt  
27 a Bayesian approach to constrain the model output with instrumental sur-  
28 face temperature and ocean heat observations. Our approach accounts for  
29 the uncertainties in the properties of model-data residuals. We use a Markov  
30 chain Monte Carlo method to obtain a posterior pdf of these parameters.

31 The mode of the climate sensitivity estimate is 2.8 °C, with the correspond-  
32 ing 95% credible interval ranging from 1.8 to 4.9 °C. These results are gen-  
33 erally consistent with previous studies. The  $CS$  pdf is sensitive to the assump-  
34 tions about the priors, to the effects of anthropogenic sulfate aerosols, and  
35 to the background vertical ocean diffusivity. Our method can be used with  
36 more complex climate models.

## 1. Introduction

37 Climate hindcasts and projections are strongly affected by two key climate model pa-  
38 rameters: climate sensitivity ( $CS$ ) and vertical ocean diffusivity. Meridional overturning  
39 circulation, global temperature, and ocean heat accumulation that produces thermosteric  
40 sea level rise are good examples of climate variables that depend on these parameters  
41 [*Goes et al.*, 2010; *Knutti et al.*, 2002]. Better characterization of the uncertainty in these  
42 parameters is thus critical for future climate prediction.

43 Climate sensitivity is defined as the equilibrium near-surface temperature response to  
44 a doubling of atmospheric  $CO_2$ .  $CS$  is a measure of climate feedbacks that amplify or  
45 dampen the direct response of near-surface temperature to radiative forcings [*Andronova*  
46 *et al.*, 2007]. Vertical ocean diffusivity is a parameter that influences heat uptake by the  
47 ocean. It parametrizes mixing processes below the grid scale of climate models. For the  
48 same climate sensitivity, at higher diffusivities the atmosphere will reach the equilibrium  
49 temperature specified by  $CS$  more slowly, due to more heat flux into the deep ocean [*NAS*,  
50 1979].

51 In order to estimate these parameters from climate models and observations, one needs  
52 to know past climate forcings. Both parameter estimation studies and simple theoretical  
53 considerations show that assumptions about these forcings influence climate sensitivity  
54 estimates and the uncertainty surrounding them [*Andreae et al.*, 2005; *Tanaka et al.*,  
55 2009; *Urban and Keller*, 2010]. For example, *Andreae et al.* [2005] use a zero-dimensional  
56 climate model to illustrate that when they assume no aerosol effects, a climate sensitivity  
57 of just  $1.3\text{ }^\circ\text{C}$  is needed to explain the observed 1940-2000 warming. On the other hand,

58 aerosol forcing of  $-1.7 \text{ W m}^{-2}$  (a value that is within the IPCC range [*Forster et al.*, 2007])  
59 requires a climate sensitivity of more than  $10 \text{ }^\circ\text{C}$  [*Andreae et al.*, 2005]. Out of the main  
60 climate forcings, the forcings due to aerosols are especially uncertain. A large part of this  
61 uncertainty is due to anthropogenic sulfate aerosols [*Forster et al.*, 2007].

62 Parameters controlling climate sensitivity, vertical diffusion in the ocean, and strength of  
63 anthropogenic sulfate aerosols are commonly estimated using model runs and observations  
64 [*Knutti et al.*, 2002, 2003; *Forest et al.*, 2002, 2006; *Drignei et al.*, 2008; *Tomassini et al.*,  
65 2007; *Edwards et al.*, 2007; *Sanso and Forest*, 2009]. Typically, an ensemble of model  
66 runs is used where the key parameters are systematically varied. The outputs from these  
67 runs are then compared with the observations, and the posterior probability distribution  
68 functions (pdfs) for model parameters are derived.

69 One conceptually simple methodology selects only the model runs that are consistent  
70 with the observations using a broad, heuristic approach [*Knutti et al.*, 2003]. In this  
71 framework all parameter combinations that pass the consistency criterion are assigned  
72 a uniform probability, while those that do not pass it receive a zero probability. These  
73 probabilities are then used to construct the posterior pdfs.

74 A more complex approach uses Bayesian statistics. This approach requires: (i) a model  
75 ensemble, (ii) observations, (iii) a statistical model that relates climate model output to  
76 the observations, and (iv) prior information about the model parameters (priors). In this  
77 framework, each parameter combination is associated with a likelihood that depends on  
78 how well the corresponding model output matches the observations [*Tomassini et al.*,  
79 2007; *Sanso and Forest*, 2009]. The likelihood,  $L(Y|\Theta)$ , describes the degree of belief that  
80 the physical observations  $Y$  came from a climate model and a statistical model (describing

81 the properties of data-model residuals) with unknown parameters  $\Theta$ . Once the statistical  
82 model is defined, the likelihood  $L(Y|\Theta)$  can be calculated from the residuals between the  
83 model output and the observations. In the Bayesian framework, the posterior probability  
84 of the unknown parameters given the observations is proportional to  $L(Y|\Theta)$ , and to the  
85 prior probability of the parameters:

$$p(\Theta|Y) \propto L(Y|\Theta) \times p(\Theta). \quad (1)$$

86 While the posterior probability  $p(\Theta|Y)$  can be evaluated on a grid of parameter values,  
87 this can become too computationally expensive if the parameter space is multidimen-  
88 sional. In such cases Markov Chain Monte Carlo (MCMC) methods [*Metropolis et al.*,  
89 1953; *Hastings*, 1970] can be used. The MCMC generates a sequence of parameter values  
90 (a Markov chain) which may be treated approximately as samples from the posterior dis-  
91 tribution. Hence, virtually any property of the posterior distribution can be approximated  
92 by a corresponding sample property of this sequence.

93 Intermediate Complexity Earth System models are frequently used for this analysis  
94 [*Forest et al.*, 2002, 2006; *Knutti et al.*, 2003; *Tomassini et al.*, 2007; *Sanso and Forest*,  
95 2009]. The appeal of these models is that they can be run at many parameter settings  
96 with relative ease. At the same time these models still represent many relevant physical  
97 processes. While the models can be run hundreds of times, many more runs at arbitrary  
98 parameter values are needed for the MCMC sampling. To approximate model output at  
99 these values, emulators (statistical approximators of climate models) can be used [e.g.,  
100 *Drignei et al.* [2008]; *Holden et al.* [2010]; *Edwards et al.* [2010]]. The emulators draw on

101 information about model outputs at known parameter settings to interpolate the output  
102 to any desired parameter setting.

103 In this study, we use the University of Victoria Earth System Climate Model (UVic  
104 ESCM) to estimate these important climate parameters. We constrain the ensemble of  
105 model runs with atmospheric surface temperature and ocean heat content observations  
106 to present probability distribution functions for key model parameters controlling the  
107 processes described above: climate sensitivity  $CS$ , background vertical ocean diffusivity  
108  $K_{bg}$ , and a scaling parameter for the direct effects of anthropogenic sulfate aerosols  $A_{sc}$ .  
109 The use of the full 3D ocean allows for the representation of the non-linear effects of  $K_{bg}$   
110 on ocean dynamics and currents (e.g., on the Meridional Overturning Circulation). We  
111 present posterior joint and marginal pdfs for the parameters, and explore the sensitivity  
112 of the results to prior assumptions.

## 2. Earth System Model, its Emulator, and Observational Constraints

### 2.1. Model Description

113 We use the University of Victoria Earth System Climate Model (UVic ESCM) [*Weaver*  
114 *et al.*, 2001] for our analysis. The atmospheric component is a one-layer energy-moisture  
115 balance model, with winds prescribed using the NCAR/NCEP climatology. The oceanic  
116 component is a three-dimensional model MOM2 [*Pacanowski*, 1995]. Both the atmo-  
117 spheric and the oceanic components have horizontal resolution of  $1.8^\circ$  (lat)  $\times$   $3.6^\circ$  (lon).  
118 The ocean has 19 depth levels. The model also includes terrestrial vegetation and carbon  
119 cycle [*Cox*, 2001], oceanic biogeochemistry based on *Schmittner et al.* [2005], and ther-  
120 modynamic sea ice. We use the modified 2.8 version of the model. Specifically, we use a

121 newer solar forcing, and include new transient forcings. The new forcings are described  
 122 in Section 2.3.

## 2.2. Model parameters

### 123 2.2.1. Climate Sensitivity ( $CS$ )

124 Climate sensitivity is defined as the equilibrium response of global average near-surface  
 125 temperature to a doubling of atmospheric  $CO_2$ . Climate sensitivity is a diagnosed param-  
 126 eter in the UVic ESCM. We vary  $CS$  through an additional parameter  $f^*$  that perturbs  
 127 local outgoing longwave radiation:

$$Q_{PLW}^* = Q_{PLW} + f^*(T_t - T_0). \quad (2)$$

128 Here  $T_o$  is temperature at equilibrium (i.e. at the start of the transient run),  $T_t$  is a  
 129 temperature at time  $t$ ,  $Q_{PLW}$  is the planetary outgoing longwave radiation as calculated  
 130 in the standard 2.8 version of the model and  $Q_{PLW}^*$  represents the modified outgoing  
 131 longwave radiation. This approach is similar to *Matthews and Caldeira* [2007] and *Zickfeld*  
 132 *et al.* [2009], but here the temperature terms are functions of latitude and longitude.

133 While  $f^*$  is the input parameter to the model, we want to know the  $CS$  values for each  
 134 ensemble model run (Section 2.3). We determine the relationship between  $f^*$  and  $CS$   
 135 using a small number of  $CO_2$  doubling experiments with varying  $f^*$  values at  $K_{bg} = 0.1$   
 136  $cm^2 s^{-1}$ . The runs continue for 2250 years to capture the equilibrium response of the  
 137 model to  $CO_2$ . The  $CS$  is diagnosed as the average global temperature during the last 50  
 138 years of the runs minus the 50 year average prior to doubling. This mapping neglects a

139 potential dependency of  $CS$  on  $K_{bg}$  at the same value of  $f^*$ . We adopt a prior range for  
 140  $CS$  from 1.1 to 11.2 (Table 1).

### 141 **2.2.2. Background Vertical Ocean Diffusivity ( $K_{bg}$ )**

142 The rate at which surface temperatures adjust to radiative forcings is controlled by  
 143 the rate at which heat is absorbed by the ocean. The vertical mixing of heat in the  
 144 ocean is parameterized in UVic ESCM by a vertical diffusivity parameter  $K_v$ , which has  
 145 contributions from tidal and background diffusivities [*Schmittner et al.*, 2009]:

$$K_v = K_{tidal} + K_{bg}. \quad (3)$$

146  $K_{tidal}$  uses the parameterization of *St. Laurent et al.* [2002] following the methodology  
 147 of *Simmons et al.* [2004]. The background diffusivity  $K_{bg}$  is assumed to be globally  
 148 uniform. We vary  $K_{bg}$  to obtain different vertical ocean diffusivities ( $K_v$ ), while keeping  
 149 standard parameters for  $K_{tidal}$ . In our model,  $K_{bg}$  largely determines the total diffusivity  
 150 in most areas of the pelagic pycnocline since the tidal component is small in those areas  
 151 [*St. Laurent et al.*, 2002; *Schmittner et al.*, 2009]. As in *Schmittner et al.* [2009] and *Goes*  
 152 *et al.* [2010], the model is modified to limit  $K_v$  to  $\geq 1 \text{ cm}^2 \text{ s}^{-1}$  in the Southern Ocean  
 153 below 500 m for better agreement with observations. Following *Goes et al.* [2010], we  
 154 adopt the prior range for  $K_{bg}$  from 0.1 to 0.5  $\text{cm}^2 \text{ s}^{-1}$  (Table 1).

### 155 **2.2.3. Anthropogenic Aerosol Scaling Factor ( $A_{sc}$ )**



Direct anthropogenic sulfate effects are modeled through spatially-resolved sulfate albedo  $\Delta a_s$  following *Matthews et al.* [2004] and *Charlson et al.* [1991] according to:

$$\Delta a_s = A_{scl} \frac{\beta \tau (1 - \alpha_s)^2}{\cos(Z_{eff})}. \quad (4)$$

156 Here  $\beta = 0.29$  is the upward scattering parameter,  $\tau$  is the aerosol optical depth field,  $\alpha_s$  is  
 157 surface albedo, and  $Z_{eff}$  is the effective solar zenith angle. The strength of anthropogenic  
 158 sulfate aerosol effects is modulated via the scaling parameter ( $A_{sc}$ ). This parameterization  
 159 does not account for the indirect effects of the sulfates on clouds. However, the indirect  
 160 effects were found to be roughly proportional to the direct effects on major components  
 161 of the Earth's radiation budget and climate on the global scale under idealized climate in  
 162 a study by *Bauer et al.* [2008]. We use the prior range for  $A_{sc}$  from 0 to 3 (Table 1).

### 2.3. Hindcast Model Runs

163 We run an ensemble of UVic ESCM model runs where we systematically vary the three  
 164 parameters over their prior ranges. Specifically,  $K_{bg}$  is varied on a uniform grid with values  
 165 of (0.1, 0.2, 0.3, 0.4, 0.5)  $\text{cm}^2 \text{s}^{-1}$ . We sample  $CS$  at (1.14, 1.64, 2.15, 2.62, 3.11, 3.98,  
 166 5.36, 6.51, 8.20, 11.2)  $^{\circ}\text{C}$ . The samples for  $A_{sc}$  are (0, 0.75, 1.5, 2.25, 3). These values  
 167 form a quasi-cubic grid (Figure 4).

168 We spin the model up from observed data fields for 3,500 years with forcings set at year  
 169 1800 values. The transient runs continue from year 1800 to the present using historic  
 170 radiative forcings. Volcanic aerosols, anthropogenic sulfate aerosols, changes in solar  
 171 constant, and additional greenhouse gases such as  $\text{CH}_4$ ,  $\text{N}_2\text{O}$  and CFCs, are implemented  
 172 following *Goes et al.* [2010]. Specifically, the volcanic radiative forcing anomalies are from  
 173 *Crowley* [2000a, b] for the period from 1800-1850, and from *GISS* [2007] and *Sato et al.*

174 [1993] for years 1850 to 2000. We update the solar forcing using the data of *Krivova et al.*  
175 [2007]. The atmospheric CO<sub>2</sub> concentration forcing is from *Etheridge et al.* [1998] and  
176 *Keeling et al.* [2004], complemented by the RCP8.5 scenario data after year 2002 [*Moss*  
177 *et al.*, 2010; *Riahi et al.*, 2007].

## 2.4. Observational Constraints

178 We use two observational constraints. The first is global average atmospheric surface  
179 / ocean surface temperatures ( $T$ ) from the HadCRUT3 dataset of the Hadley Center  
180 [*Brohan et al.*, 2006]. These observations are defined as anomalies with respect to the  
181 1850-1899 period average. The observations cover the time period from 1850 to 2006  
182 (Figure 2). The second constraint is global total ocean heat content ( $OHC$ ) in the 0-700  
183 m layer [*Domingues et al.*, 2008]. These observations span the period from 1950 to 2003,  
184 and are calculated as anomalies with respect to the whole observation period (Figure 2).  
185 Modeled temperature and ocean heat content are converted to anomalies to be consistent  
186 with the observational constraints.

## 2.5. Gaussian Process Emulator

187 The MCMC sampling requires a large number of model runs ( $> 10000$ ) at arbitrary  
188 parameter values. Since it is computationally infeasible to run UVic ESCM at that many  
189 parameter settings, we use a statistical emulator that can approximate the model outputs  
190 at any parameter value. We adopt Gaussian Process ( $GP$ ) emulation. This technique  
191 was previously used to approximate climate models by *Bhat* [2010], *Sanso and Forest*  
192 [2009] and *Rougier et al.* [2009]. We emulate model output as a function of climate  
193 parameters separately for temperature and for ocean heat content. For each tracer, we

194 develop separate emulators for each time step during the years for which the observations  
 195 are available (Section 2.4). Thus, we build a total of 157 emulators for temperature, and  
 196 54 for the ocean heat content.

197 We define model output of tracer  $k$  at time  $t$  as  $f_{t,k}(\theta)$  where  $\theta$  is a vector of model  
 198 parameters  $(K_{bg}, CS, A_{sc})$ . The  $f_{t,k}(\theta)$  is only defined on a discrete set of parameter values  
 199 where the model was run. The purpose of the emulator is to estimate a function  $\tilde{f}_{t,k}(\theta)$   
 200 approximating model output on the continuous parameter ranges specified in Table 1. In  
 201 the following discussion we will consider the emulator for atmospheric surface temperature  
 202 at time  $t_0$ . The emulators for all other times and for the second tracer (ocean heat content)  
 203 follow a similar statistical model. The indices  $t$  and  $k$  will thus be dropped from the rest  
 204 of the emulator description.

205 The emulator is developed in linearly rescaled coordinates with transformed parameters  
 206  $\theta' = (K'_{bg}, CS', A'_{sc})$  each taking on a range from zero to unity. The emulator approximates  
 207 the climate model output as:

$$\tilde{f}(\theta') = P(\theta') + Z(\theta'), \quad (5)$$

208 where  $P$  is a quadratic polynomial in model parameters, and  $Z$  is a zero-mean Gaussian  
 209 Process with an isotropic covariance function. Specifically, the covariance between  $Z$  at  
 210 parameters  $\theta'_i$  and  $\theta'_j$  is modeled as  $mC(i, j)$  where  $m$  is a scale multiplier and  $C$  is defined  
 211 by:

$$C(i, j) = \exp \frac{-D_{ij}}{l}. \quad (6)$$

212 Here  $D_{i,j}$  is the Euclidean distance between the two model parameter settings and  $l$  is a  
213 range parameter. Based on exploratory data analysis, we choose  $l=0.6$ . This formulation  
214 ensures that model output at nearby parameter settings is highly correlated (i.e. model  
215 output is a smooth function of the parameters). We choose a nugget variance  $\sigma_\epsilon^2$  of zero.  
216 This implies that the emulator is equal to model output at the points of the ensemble  
217 design grid.

218 We estimate the polynomial parameters and  $m$ . The polynomial parameters are the  
219 generalized linear squares estimates adjusting for the covariance function of the spatial  
220 process. They have a closed form solution that follows a standard formulation in Universal  
221 Kriging.  $m$  is likewise found by maximum likelihood given the parameter  $\lambda = \sigma_\epsilon^2/m = 0$ ,  
222 and it has a closed form solution given  $\lambda$  as well (D. Nychka, personal communication).  
223 The optimized parameters provide the Best Linear Unbiased Estimate (BLUE) of  $\tilde{f}(\theta')$   
224 [Furrer *et al.*, 2010].

225 Emulators for other times and variables have different  $P$  and  $m$ . Henceforth all the  
226 emulators for all time steps and both tracers will be collectively referred to as the “emu-  
227 lator”.

228 The emulator was extensively tested using the leave-one-out cross validation analysis.  
229 The emulator is found to perform adequately well (e.g., Figure 1) during the times when  
230 the variability of model output across the parameter space is high. The cross-validation  
231 errors are larger in the relative sense during the times close to the midpoints for the  
232 averaging periods for the anomalies (i.e. year 1870 for temperature, and 1980 for ocean  
233 heat content). At such times the signal is small and the model output is not a smooth  
234 function of the parameters, therefore it is impossible to accurately predict it based on the

235 information from the remaining runs. We are unaware of any improvement in emulation  
 236 techniques that could overcome this problem. We note that in this case the emulator  
 237 errors are very low in the absolute sense and they are not expected to affect the estimation  
 238 results. Overall, based on the cross-validation analysis, we are confident that the emulator  
 239 provides a reasonable tool to interpolate model output.

### 3. Statistical Model and Markov Chain Monte Carlo

240 We use a Bayesian parameter estimation method. In order to be able to evaluate the  
 241 likelihood of observations given the unknown parameters  $L(Y|\Theta)$ , we need a statistical  
 242 model that defines the relationship between the model (and the emulator) output and the  
 243 observations. We refer to the emulator output by  $\tilde{f}_{t,k}(\theta)$  (for time  $t$ , tracer  $k$ , and param-  
 244 eter combination  $\theta$ ). The observations are denoted by  $y_{t,k}$ . We denote each observational  
 245 time series by  $\mathbf{y}_k = y_{1,k}, \dots, y_{N_k,k}$  where  $N_k$  is the number of observations for tracer  $k$ . The  
 246 set of all observations is referred to as  $Y = (\mathbf{y}_T, \mathbf{y}_{OHC})$ .

247 We assume that the discrepancy between the emulator and the observations is due to  
 248 the time constant bias  $b_k$  and time-varying error  $\epsilon_{t,k}$ . Thus, our statistical model is:

$$y_{t,k} = \tilde{f}_{t,k} + b_k + \epsilon_{t,k}. \quad (7)$$

249  $\epsilon_{t,k}$  results from (i) model error, (ii) natural climate variability, (iii) emulator error,  
 250 and (iv) observational error. We assume that  $\epsilon_{t,k}$  is an autoregressive process of order  
 251 1 (AR1) with unknown AR1 parameters  $\sigma_k^2$  and  $\rho_k$ .  $\sigma_k^2$  represents the variance of the  
 252 AR(1) innovations while  $\rho_k$  represents the autocorrelation of lag1 (i.e. correlation of  
 253  $\epsilon_{t,k}$  with  $\epsilon_{t-1,k}$ ). This form is chosen both for its simplicity and the ability to account

254 for the uncertain autocorrelation in the error terms. The bias term  $b_k$  represents time-  
 255 independent biases. Note that for ocean heat content we use anomalies with respect to  
 256 the entire observational period. As a result, the average modeled and observed  $OHC$  is 0  
 257 by definition and we set  $b_{OHC}$  to 0. Our statistical model is similar to *Urban and Keller*  
 258 [2010], although they do not incorporate bias terms.

For this statistical model, the likelihood of each observational time series  $\mathbf{y}_k$  given the  
 UVic ESCM model output and the statistical parameters  $L(\mathbf{y}_k|\theta, \sigma_k, b_k, \rho_k)$  is given by  
 [Bence, 1995] and is provided in the Appendix. We assume independence between the  
 model-data residuals for different tracers. Under this assumption, the likelihood of both  
 observations is equal to the product of the individual likelihoods:  $L(Y) = L(\mathbf{y}_1) \times L(\mathbf{y}_2)$ .  
 Denote the set of all parameters by  $\Theta = (K_{bg}, CS, A_{sc}, \sigma_T, \rho_T, b_T, \sigma_{OHC}, \rho_{OHC})$ . Using  
 Bayes Theorem, the posterior probability of the parameters can be calculated as:

$$p(\Theta|Y) \propto L(Y|\Theta) \times p(\Theta) \quad (8)$$

259 where  $p(\Theta)$  is the prior for the parameters (Section 4).

260 Two distinct approaches to estimate the properties of the the error process  $\epsilon$  are (i) from  
 261 the observations or models [Forest et al., 2006; Tomassini et al., 2007], or (ii) directly from  
 262 the model-data residuals together with the physical parameters [Urban and Keller, 2010;  
 263 Goes et al., 2010; Tonkonojenkova, 2010]. Here we use the second approach and estimate  
 264 all parameters together in the MCMC step.

265 We draw samples from the joint posterior  $p(\Theta|Y)$  using the MCMC algorithm [Metropo-  
 266 lis et al., 1953; Hastings, 1970] and generate the posterior probability distribution of  $\Theta$ .  
 267 Our MCMC prechains are 50,000 members long, while the final chain has 300,000 mem-  
 268 bers. We use information from previous chain covariance to construct the proposal dis-

269 tribution for each subsequent chain following *Roberts and Rosenthal* [2009]. We test the  
270 chains for convergence using the MCMC standard errors from the consistent batch means  
271 procedure [*Flegal et al.*, 2008; *Jones et al.*, 2006], and by repeating the assimilation with  
272 different starting values of the parameters for the final chain. Neither of these checks sug-  
273 gest any issues with convergence. Hence, we are satisfied that our MCMC-based inference  
274 provides reasonable estimates of the posterior pdfs.

#### 4. Priors

275 We run two assimilation experiments. In the base case experiment we use non-uniform  
276 priors for climate sensitivity and background vertical ocean diffusivity. We refer to this  
277 experiment as NON-UNIF. The priors for this experiment are listed in Table 1 and plotted  
278 in Figure 3. For  $K_{bg}$  the prior is Lognormal  $(-1.55, 0.59)$   $\text{cm}^2 \text{s}^{-1}$  [*Bhat*, 2010]. This prior  
279 has a mode of  $0.15 \text{ cm}^2 \text{s}^{-1}$  and a mean of  $0.24 \text{ cm}^2 \text{s}^{-1}$ . The prior represents our prior belief  
280 that the values of  $0.1 - 0.2 \text{ cm}^2 \text{s}^{-1}$  are more likely than  $0.4 - 0.5 \text{ cm}^2 \text{s}^{-1}$  which is suggested  
281 by *Goes et al.* [2010] who use vertical oceanic tracer distributions to constrain  $K_{bg}$ . The  
282 climate sensitivity prior incorporates weak prior information derived from current mean  
283 climate and Last Glacial Maximum constraints. Specifically, we use a product of normal  
284 inverse Gaussian distributions (*NIG*) of  $NIG(\alpha = 1.8, \delta = 2.3, \beta = 1.2, \mu = 1.7)$  and  
285  $NIG(\alpha = 1.9, \delta = 3.3, \beta = 1.0, \mu = 1.3)$ . We choose these distributions for their empirical  
286 ability to simultaneously fit the lower, upper, and best estimates in *Knutti and Hegerl*  
287 [2008], not because we have any theoretical motivation for the *NIG* distribution. While  
288 the central tendencies of the two *NIG* pdfs are generally compatible with past studies, the  
289 distributions are not based on any specific pdf from any of these studies. The combined  
290 prior distribution for  $CS$  is shown in Figure 3. It has a mean of  $3.25 \text{ }^\circ\text{C}$ , and the 90%

291 interval from 1.7 to 5.2 °C. We use uniform priors for  $A_{sc}$  and for all statistical parameters  
292 over the ranges specified in Table 1.

293 To explore the sensitivity of the results to priors, we run a second assimilation exper-  
294 iment, where all priors are uniform over the ranges shown in Table 1. We refer to this  
295 experiment as UNIF.

## 5. Results

### 5.1. Probabilistic Hindcasts

296 The probabilistic hindcasts capture the overall temporal structure of the observations  
297 (Figure 2). Specifically, the emulator is able to correctly represent the trend due to  
298 greenhouse warming (black line). We add an AR1 error process (representing model,  
299 observational, and emulator error, as well as the natural variability) to each emulator  
300 from the sub-sampled MCMC chain to produce the 95% credible intervals. In case of  
301 temperature, each emulator is corrected by adding a corresponding bias term from the  
302 chain. Overall, the method produces a reasonable surprise index (e.g., 1.9% of the ocean  
303 heat content and 5.1% of the temperature observations lie outside of the 95% hindcast  
304 range).

305 The surface air temperature from the best fit emulator illustrates the effects of the  
306 stratospheric volcanic aerosols, with several prominent short-term coolings associated with  
307 the eruptions. For some of these eruptions, such as Agung (1963) and Mount Pinatubo  
308 (1991), the modeled response matches the observations relatively well, while for others,  
309 such as Krakatoa (1883), the model displays considerable cooling that is not present in the  
310 observations. Some of this discrepancy might be due to the unresolved climate variability,



311 and due to the uncertainty in the past volcanic radiative effects [*Ammann et al.*, 2003]  
312 and temperature observations.

## 5.2. Parameter Estimates

313 Under the baseline assumptions of non-uniform priors, posterior pdfs for climate sen-  
314 sitivity and vertical ocean diffusivity are broadly consistent with previous studies. The  
315 mode of the climate sensitivity pdf is 2.8 °C, and the mean is 3.1 °C. The 95% posterior  
316 credible interval ranges from 1.8 °C to 4.9 °C (Table 2). These values are broadly consis-  
317 tent with the likely range of 2 to 4.5 °C, and the most likely value of 3 °C given by the  
318 IPCC [*Solomon et al.*, 2007]. The mode is similar to results from *Forest et al.* [2006] and  
319 *Knutti et al.* [2003], and is slightly higher than in *Tomassini et al.* [2007].

320 For  $K_{bg}$ , we estimate a mode of 0.11 cm<sup>2</sup> s<sup>-1</sup>, and a mean of 0.19 cm<sup>2</sup> s<sup>-1</sup> (Table 2, Figure  
321 3). The pdf for  $K_{bg}$  was reported to depend on the tracers used to constrain this parameter  
322 [*Schmittner et al.*, 2009]. The mode of the  $K_{bg}$  matches results of *Schmittner et al.* [2009]  
323 based on global vertical ocean profiles of CFC11, and of  $\Delta^{14}C$ , and is slightly lower than  
324 0.15 cm<sup>2</sup> s<sup>-1</sup> reported in [*Goes et al.*, 2010] based on profiles of three tracers. We stress  
325 that  $K_{bg}$  is not directly comparable with vertical diffusivities in other models [*Tomassini*  
326 *et al.*, 2007; *Kriegler*, 2005] because these parameters represent different processes. For  
327 example, our  $K_{bg}$  excludes tidally induced and Southern Ocean mixing, while the related  
328  $K_v$  in *Kriegler* [2005] accounts for all vertical mixing processes. Therefore, our results  
329 should be interpreted as specific to our version of UVic ESCM.

330 The estimated aerosol scaling factor has the most likely value of 1.2. This is broadly  
331 consistent with the default assumptions on the aerosol effects in the UVic ESCM (which  
332 imply the value of 1). Estimation of  $A_{sc}$  should be interpreted with caution because  $A_{sc}$

333 implicitly includes effects due to neglected forcings that might have emission or concen-  
334 trations patterns similar to the anthropogenic sulfates. To better constrain  $A_{sc}$  it will be  
335 necessary to include these neglected forcings. Otherwise, one could interpret the value  
336 of  $A_{sc}$  as representing the combined effects of the aerosols as well as the neglected forc-  
337 ings. Similar to the case of  $K_{bg}$ ,  $A_{sc}$  is a model specific parameter and can not be readily  
338 compared to results from other models [i.e. *Tanaka et al.* [2009]].

339 As in previous studies, the climate sensitivity pdf, and its upper tail in particular,  
340 are sensitive to the assumptions about the priors [e. g. *Forest et al.*, 2002, 2006; *Sanso*  
341 *and Forest*, 2009; *Tomassini et al.*, 2007; *Annan and Hargreaves*, 2011] (Figure 3). For  
342 example, replacing the expert prior with the uniform prior moves the upper bound of  
343 the 95% credible interval for  $CS$  to 10.2 °C (Table 2). This is in agreement with the  
344 results from *Forest et al.* [2006], but considerably higher than in *Annan and Hargreaves*  
345 [2011]. This discrepancy might be at least in part because *Annan and Hargreaves* [2011]  
346 consider a different type of constraint - Earth Radiation Budget Experiment (ERBE) data  
347 analyzed by *Forster and Gregory* [2006]. For the uniform prior, there is a considerable  
348 probability mass above the upper bound of the IPCC likely range of 4.5 °C (Figure 3),  
349 similar to previous studies (e.g., *Forest et al.* [2006]; *Knutti et al.* [2003]).

350 The use of uniform priors for climate sensitivity can be problematic as the posterior  
351 estimates are sensitive to the upper bound for the prior [*Annan and Hargreaves*, 2011].  
352 In addition, such priors do not take independently collected evidence from other studies  
353 into account. High climate sensitivities become possible in this case because the flat prior  
354 assigns them high weight to begin with, while the constraint provided by the observations

355 can be relatively weak. This suggests that it is crucial to use independent prior information  
356 during  $CS$  estimation whenever possible.

357 In addition, in the UNIF experiment the posterior pdf of  $K_{bg}$  is bimodal (Figure 3). Mul-  
358 timodal pdfs for  $K_{bg}$  have been previously reported by *Forest et al.* [2002] and *Tomassini*  
359 *et al.* [2007]. It is, thus far, unclear which physical mechanisms, if any, are driving this  
360 bimodality. Note that here we withhold information on vertical tracer distributions that  
361 is needed to constrain  $K_{bg}$  and that the bimodality essentially disappears once that con-  
362 straint is introduced as a prior in the NON-UNIF case.

363 Joint bivariate pdfs for parameter pairs exhibit a complex structure (Figure 4), similar  
364 to the results from from *Tomassini et al.* [2007]. Although this is not visibly evident,  
365 there is some correlation between  $K_{bg}$  and  $CS$ . Specifically, the correlation is 0.24 in the  
366 NON-UNIF experiment, and 0.44 in the UNIF experiment. This is in agreement with 0.4  
367 found in *Urban and Keller* [2010] even though the two studies differ in terms of climate  
368 models, observational constraints, and priors. It is difficult to compare these results with  
369 other studies ([e. g. *Tomassini et al.*, 2007; *Forest et al.*, 2002, 2006]) because they do  
370 not report the numerical value for the correlation coefficient while the pairs plots of the  
371 parameters can underestimate the correlation [*Urban and Keller*, 2010].

372 Climate sensitivity is even more strongly correlated with  $A_{sc}$ , meaning that for higher  
373 climate sensitivity, higher aerosol effects are needed to explain historical climate change.  
374 This agrees with results from *Andreae et al.* [2005] and *Tanaka et al.* [2009] and implies  
375 that reducing uncertainty in  $A_{sc}$  will help reduce uncertainty in climate sensitivity. Ruling  
376 out high values of  $A_{sc}$  is especially important, because this is where climate sensitivity  
377 pdf appears to be most sensitive to  $A_{sc}$  (Figure 4).

378 When the uniform priors on  $K_{bg}$  and  $CS$  are used, higher aerosol scaling values become  
379 possible, even though the prior on  $A_{sc}$  is the same in both cases. Because  $A_{sc}$  and  $CS$   
380 are correlated, higher aerosol scalings are necessary to counteract higher warming due to  
381 larger climate sensitivities in the uniform prior case to match the observations.

382 Climate parameter estimation using a model with a 3D ocean (GENIE-1) has been  
383 previously performed by *Holden et al.* [2010] so it might be interesting to compare our  
384 methodology and results with that study. *Holden et al.* [2010] vary a much larger set  
385 of parameters and derive a pdf for climate sensitivity using a Last Glacial Maximum  
386 (LGM) tropical Sea Surface Temperature (SST) anomaly as a main constraint. They also  
387 indirectly use information from several global climate metrics through a pre-calibration  
388 procedure. In our study we consider an orthogonal set of constraints that includes infor-  
389 mation about the time-resolved response of climate to modern forcings. We also provide  
390 a probabilistic estimate of vertical ocean diffusivity  $K_{bg}$ . In terms of the ocean models  
391 used, *Holden et al.* [2010] employ a coarse resolution frictional geostrophic model. On the  
392 other hand, the resolution of UVic ESCM is much higher and the dynamics is based on  
393 the Navier-Stokes equations, subject to the hydrostatic and Boussinesq approximations.  
394 The statistical methodologies are different as well. In particular, our approach is fully  
395 Bayesian and we use explicit priors for all model parameters. Also, the statistical proper-  
396 ties of the error process are assumed in *Holden et al.* [2010], while here we estimate them  
397 together with the physical model parameters. The mode of climate sensitivity found in  
398 *Holden et al.* [2010] is 3.6 °C under the favored set of assumptions, which is substantially  
399 higher than 2.8 °C in our baseline case of non-uniform priors. We cannot attribute this

400 gap with certainty to any specific factor due to the number of differences between the  
401 studies.

## 6. Caveats

402 Our forthcoming conclusions are subject to several caveats. The first set of caveats  
403 deals with the Earth System model. Our model does not include all forcings (such as,  
404 sulfate effects on clouds or tropospheric ozone [*Forster et al.*, 2007]). The patterns of some  
405 of excluded forcings might be similar to anthropogenic sulfates, thereby biasing the  $A_{sc}$   
406 estimates. Including thus far neglected forcings is the subject of future research. Also, we  
407 only consider a subset of uncertain climate parameters. Our results would change if these  
408 additional uncertainties were considered. The model relies on a number of simplifications.  
409 The representation of open ocean mixing is highly parametrized and ignores, for example,  
410 effects of transient upper ocean mixing processes, such as tropical cyclones, that have  
411 been shown capable of influencing upper-ocean temperature patterns through mixing of  
412 heat [*Sriver et al.*, 2010]. We vary the longwave radiation feedbacks to change climate  
413 sensitivity. In reality, the uncertainty in shortwave radiative feedbacks also contributes to  
414 the  $CS$  uncertainty [*Bony et al.*, 2006]. Also, we only use a single model and neglect the  
415 uncertainty in model response to external forcings [*Stouffer et al.*, 2006]. Finally, we do  
416 not fully account for past climate forcings uncertainties.

417 The second set of caveats is concerned with observations. When a short instrumental  
418 record is used, the results of our method can be influenced by natural climate variability  
419 and by observational errors comprising the residuals between the model and observations  
420 [*Tonkonojenkova*, 2010]. Adding more observations can improve the parameter estimates,  
421 as could using spatially resolved information.

422 Finally, limitations of the parameter estimation method deserve mentioning. We use  
423 a simplified likelihood function that does not account for the spectral complexity of the  
424 residuals, nor for the decrease of observational errors with time. Incorporating a more  
425 comprehensive likelihood function that captures a cross-correlation between the residuals  
426 for different tracers is the subject of future research.

## 7. Conclusions

427 Using a Bayesian approach, we fuse the UVic ESCM model with global observations  
428 to estimate background vertical ocean diffusivity ( $K_{bg}$ ), climate sensitivity ( $CS$ ), and the  
429 scaling parameter for the effects of anthropogenic sulfate aerosols ( $A_{sc}$ ). Our methodology  
430 incorporates the effects of  $K_{bg}$  on 3D ocean dynamics. We use a Gaussian Process emulator  
431 to provide a fast surrogate for the climate model at arbitrary parameter combinations.  
432 The parameter estimates can be used to make climate projections using the UVic ESCM  
433 in future studies.

434 The mode for  $K_{bg}$  is similar to previous results obtained using oceanic tracers such as  
435 CFC11, temperature, and  $\Delta^{14}C$  as constraints. The  $K_{bg}$  pdf is sensitive to the assumptions  
436 about the priors. If a uniform prior is used, then the results appear to show a bimodality,  
437 which is a potentially important result that might need further investigation.

438 Under the default assumptions of informative priors, the mode of climate sensitivity is  
439 2.8 °C, with the 95% credible interval from 1.8 °C to 4.9 °C. This mode is consistent with  
440 many previous studies but lower than reported in *Holden et al.* [2010] who also use a 3D  
441 ocean model. As in previous studies, the upper tail of the  $CS$  pdf is sensitive to priors.  
442 The  $CS$  pdf depends critically on  $A_{sc}$ , with much higher climate sensitivities likely at high

443 values of  $A_{sc}$ . The agreement with previous studies that use simpler climate models gives  
 444 more confidence to using these models to estimate climate sensitivity.

445

## Appendix

446 When the statistical model is defined as in Section 3, the likelihood of observational  
 447 time series  $\mathbf{y}_k$  coming from the model is given by [Bence, 1995]:

$$L(\mathbf{y}_k | \theta, \sigma_k, b_k, \rho_k) = (2\pi\sigma_{p,k}^2)^{-1/2} \exp\left(-\frac{1}{2} \frac{\epsilon_{1,k}^2}{\sigma_{p,k}^2}\right) \times \\ \times (2\pi\sigma_k^2)^{-(N_k-1)/2} \times \exp\left(-\frac{1}{2\sigma_k^2} \sum_{j=2}^{N_k} w_{j,k}^2\right).$$

448 Here  $\sigma_{p,k}^2 = \sigma_k^2/(1 - \rho_k^2)$  is stationary process variance,  $N_k$  is the number of observational  
 449 data points for tracer  $k$ , and  $w_{t,k} = \epsilon_{t,k} - \rho_k \epsilon_{t-1,k}$ ,  $t > 1$  are whitened errors.

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## Figure Captions

635 **Figure 1.** Top row: scatterplot of the temperature anomaly (with respect to the 1850-2006  
 636 mean, [K]) emulator predictions vs. actual model output values for years 1870, 1940, and 2000.  
 637 Specifically, each of the parameter combinations of the ensemble was taken out one at a time,  
 638 and the emulator was trained on the remaining 249 ensemble members. Then the emulator was  
 639 used to predict the missing value. The 1:1 line is also shown. Note that Y axis limits are different  
 640 for each subplot. Bottom row: same for the ocean heat content anomalies (with respect to the  
 641 1950-2003 mean, [ $1 \times 10^{22}J$ ]), for years 1960, 1980, and 2000. The emulator performance, of  
 642 course, will be different for other times not shown here.

643 **Figure 2.** Probabilistic model hindcasts (grey shaded area), maximum posterior probability  
 644 model output (‘best fit’, black line), and corresponding observations (red crosses) for the NON-  
 645 UNIF assimilation experiment: (a) global average atmospheric surface temperature anomaly with  
 646 respect to 1850-1899 mean [K] with corresponding observations of above surface / ocean surface  
 647 temperatures from the HadCRUT3 dataset [*Brohan et al.*, 2006]; (b) upper ocean (0-700m) heat  
 648 content anomaly with respect to 1950-2003 mean [ $1E22J$ ], and observations from *Domingues*  
 649 *et al.* [2008]. The grey area denotes the 95% credible intervals for model output taken from  
 650 a 1000-member subsampled MCMC chain, with corresponding AR1 error processes (and bias  
 651 terms for temperature) added. For the AR1 process simulations, the  $\sigma$  and  $\rho$  parameters were  
 652 taken from the corresponding chain member. For the best fit model output for temperature, the  
 653 maximum posterior probability model output was combined with the corresponding bias term.

654 **Figure 3.** Posterior pdfs (top row) and cdfs (bottom row) for model parameters obtained using  
 655 both temperature and ocean heat content observations. Red: for the NON-UNIF experiment;  
 656 blue: for the UNIF experiment. The dashed probability distribution lines represent the priors



657 used in the NON-UNIF experiment. The dashed whiskers in the box-and-whisker plots extend  
658 to the most extreme data point which is no more than 1.5 interquartile ranges from the box.

659 **Figure 4.** Bivariate joint pdfs for model parameters. Left: for the NON-UNIF experiment,  
660 right: for the UNIF experiment. The contour lines delineate the 90% and 95% posterior credible  
661 intervals. A 1000-member thinned MCMC chain is plotted using red dots. Parameters used for  
662 the UVic ESCM ensemble are shown in thick black circles.

Table 1: Prior ranges for the parameters used in the NON-UNIF experiment. Subscript  $T$  refers to the temperature data, and  $OHC$  refers to the ocean heat content data.

Parameter	Units	Lower Bound	Upper Bound	Prior Form
$K_{bg}$	$\text{cm}^2 \text{ s}^{-1}$	0.1	0.5	Lognormal(-1.55, 0.59)
$CS$	$^{\circ}\text{C}$ per $\text{CO}_2$ doubling	1.1	11.2	$NIG(\alpha = 1.8, \delta = 2.3, \beta = 1.2, \mu = 1.7) \times NIG(\alpha = 1.9, \delta = 3.3, \beta = 1.0, \mu = 1.3)$
$A_{sc}$	unitless	0	3	uniform
$\sigma_T$	$^{\circ}\text{C}$	0.01	inf	uniform
$\sigma_{OHC}$	$1 \times 10^{22} \text{ J}$	0.01	inf	uniform
$\rho_T$	unitless	0.01	0.99	uniform
$\rho_{OHC}$	unitless	0.01	0.99	uniform
$b_T$	$^{\circ}\text{C}$	-0.51	0.50	uniform

Table 2: Properties of the posterior pdfs of all estimated parameters.

Parameter	Experiment	Mode	Mean	95% credible interval
$K_{bg}$	NON-UNIF	0.11	0.19	[0.10, 0.45]
	UNIF	0.11	0.30	[0.10, 0.50]
$CS$	NON-UNIF	2.8	3.1	[1.8, 4.9]
	UNIF	3.0	4.8	[1.6, 10.2]
$A_{sc}$	NON-UNIF	1.2	1.1	[0.35, 1.5]
	UNIF	1.6	1.2	[0.25, 1.8]
$\sigma_T$	NON-UNIF	0.10	0.10	[0.091, 0.11]
	UNIF	0.10	0.10	[0.091, 0.11]
$\sigma_{OHC}$	NON-UNIF	2.6	2.7	[2.2, 3.3]
	UNIF	2.6	2.7	[2.2, 3.3]
$\rho_T$	NON-UNIF	0.58	0.58	[0.44, 0.72]
	UNIF	0.58	0.58	[0.44, 0.72]
$\rho_{OHC}$	NON-UNIF	0.079	0.17	[0.018, 0.43]
	UNIF	0.091	0.17	[0.018, 0.42]
$b_T$	NON-UNIF	-0.031	-0.031	[-0.079, 0.021]
	UNIF	-0.034	-0.033	[-0.083, 0.022]

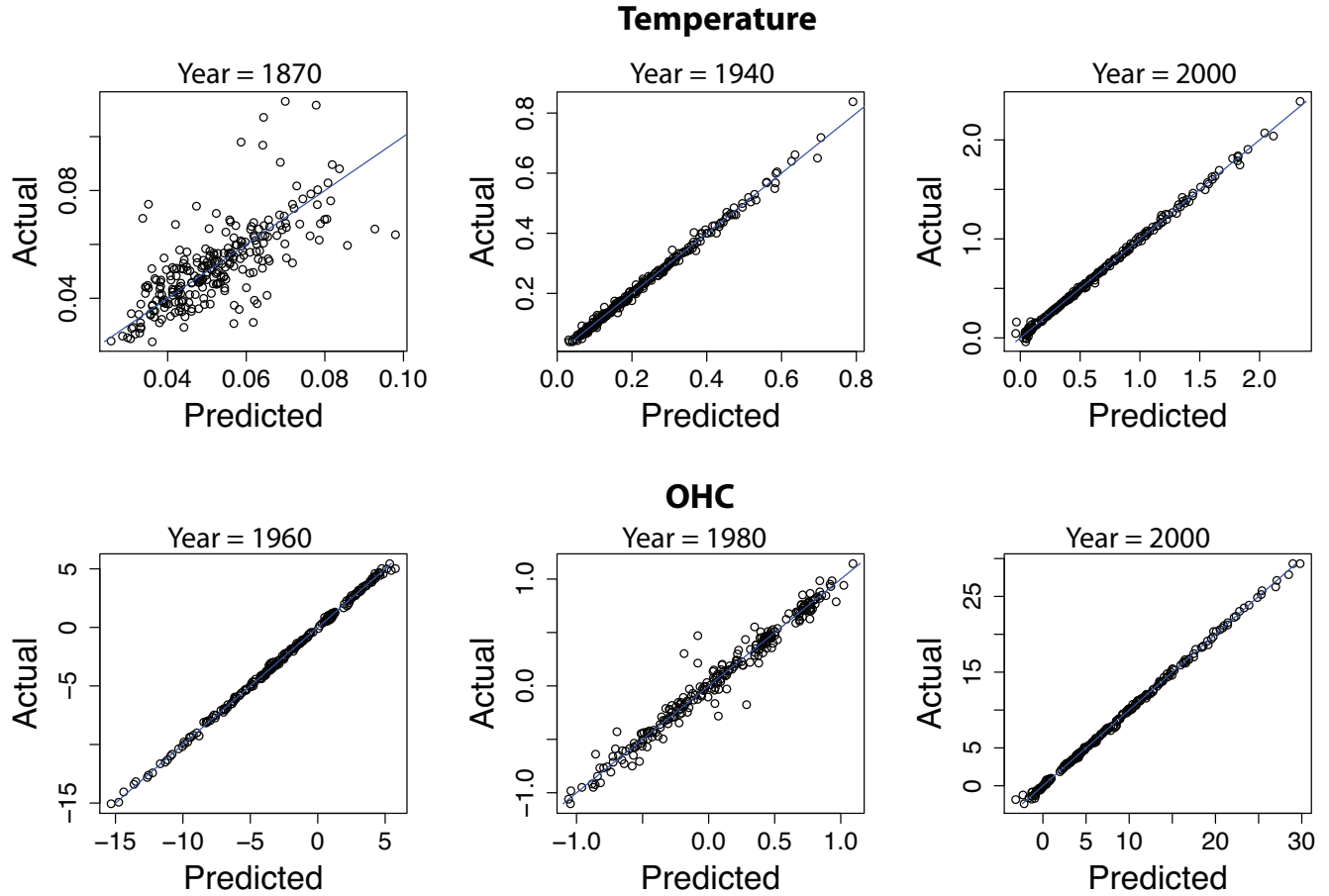


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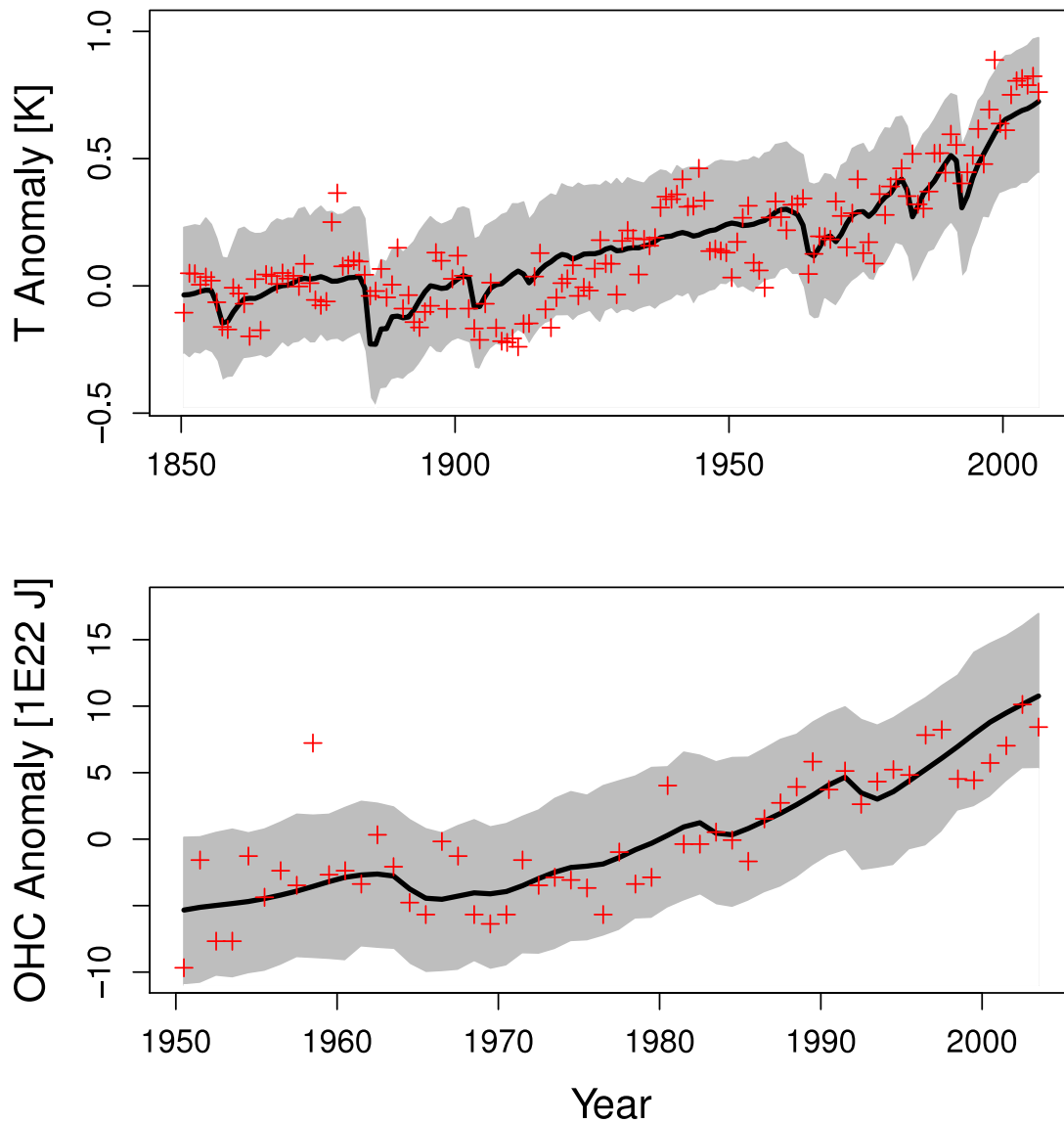


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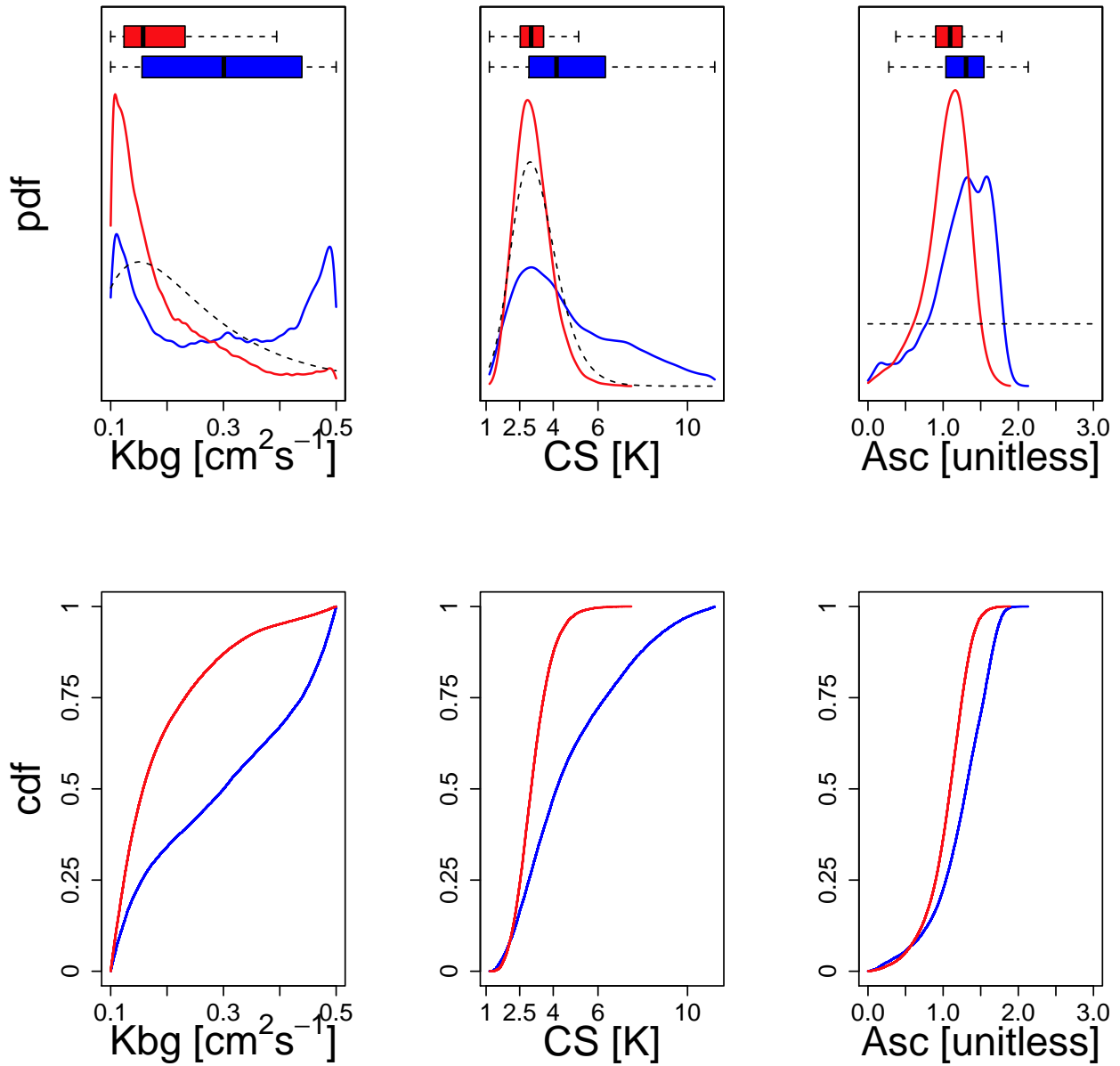


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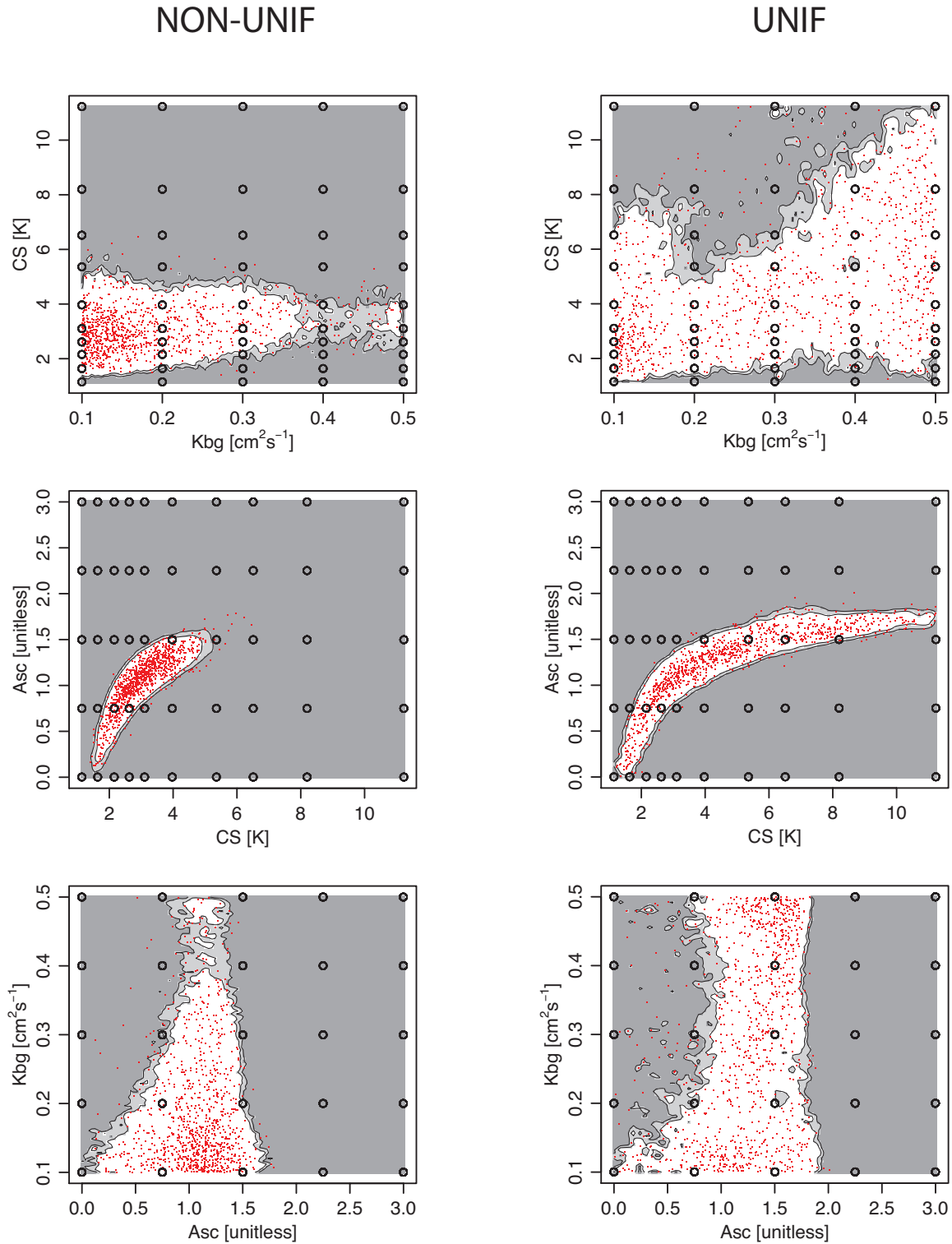


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